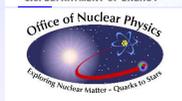


Light Quark Confinement And The Trajectory Of The Pseudoscalar Meson

STEWART V. WRIGHT

Physics Division

Argonne National Laboratory



Nonperturbative Methods in QCD

- Green function methods



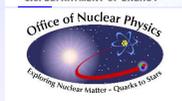
Nonperturbative Methods in QCD

- Green function methods
- Light front relativistic quantum mechanics



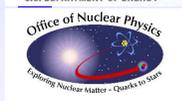
Nonperturbative Methods in QCD

- Green function methods
- Light front relativistic quantum mechanics
- Numerical studies of lattice-regularised QCD



Nonperturbative Methods in QCD

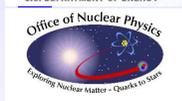
- Green function methods
- Light front relativistic quantum mechanics
- Numerical studies of lattice-regularised QCD
- QCD modelling



Nonperturbative Methods in QCD

- Green function methods
- Light front relativistic quantum mechanics
- Numerical studies of lattice-regularised QCD
- QCD modelling

Dyson–Schwinger equations



Contemporary Reviews

- **Dyson–Schwinger Equations:
Density, Temperature and Continuum Strong
QCD**

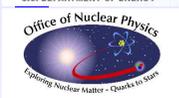
C.D. Roberts and S.M. Schmidt, nucl-th/0005064
Prog. Part. Nucl. Phys. 45 (2000) S1

- **The IR behaviour of QCD Green's functions:
Confinement, DCSB and hadrons ...**

R. Alkofer and L. von Smekal, hep-ph/0007355
Phys. Rept. 353 (2001) 281

- **Dyson–Schwinger Equations:
A Tool for Hadron Physics**

P. Maris and C.D. Roberts, nucl-th/0301049
Int. J. Mod. Phys. E12 (2003) pp. 297–365



Dyson–Schwinger Equations



Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD



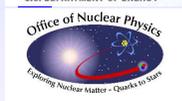
Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD
- Simplest level: Generating tool for perturbation theory



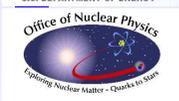
Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD
 - Hadrons are composites of Quarks and Gluons



Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD
 - Hadrons are composites of Quarks and Gluons
 - Quantitative and Qualitative importance of:
 - Dynamical Chiral Symmetry Breaking



Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD
 - Hadrons are composites of Quarks and Gluons
 - Quantitative and Qualitative importance of:
 - Dynamical Chiral Symmetry Breaking
 - Quark and Gluon confinement

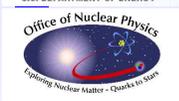


Dyson–Schwinger Equations

- Nonperturbative, continuum approach to QCD
 - Hadrons are composites of Quarks and Gluons
 - Quantitative and Qualitative importance of:
 - Dynamical Chiral Symmetry Breaking
 - Quark and Gluon confinement
- ⇒ Understanding Infrared (long range) behaviour of $\alpha_s(Q^2)$

Maris, Roberts *Phys. Rev.* **C56** 3369

Maris, Tandy *Phys. Rev.* **C60** 055214



Dressed Quark Propagator

$$S(p) = \frac{Z(p^2)}{i\gamma \cdot p + M(p^2)} \quad \longrightarrow \quad \text{---} \bullet \text{---} = \text{---} \bullet \text{---} \overset{\text{---} \bullet \text{---}}{\curvearrowright}$$

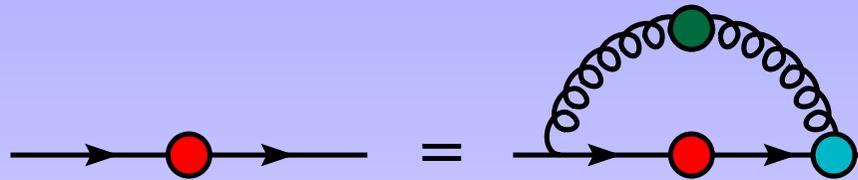
- Weak coupling expansion: Reproduces every diagram in Perturbation Theory
- Perturbation Theory:

$$M(p^2) = m_0 \left(1 - \frac{3\alpha}{4\pi} \ln \left[\frac{p^2}{m_0^2} \right] + \mathcal{O}(\alpha^2) \right)$$

$$\begin{aligned} m &\rightarrow 0 \\ &\rightarrow 0 \end{aligned}$$

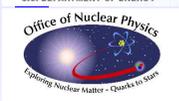
A Challenge

- Infinitely Many Coupled Equations



A Challenge

- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations



A Challenge

- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations
- Coupling between equations necessitates truncation



A Challenge



- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations
- Coupling between equations necessitates truncation
 - Weak coupling expansion
⇒ Perturbation Theory
Not useful for nonperturbative problems

A Challenge



- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations
- Coupling between equations necessitates truncation
- “Many Body” physics approach

A Challenge



- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations
- Coupling between equations necessitates truncation
- “Many Body” physics approach
 - Rainbow–Ladder (\sim Hartree–Fock) truncation. Widely explored

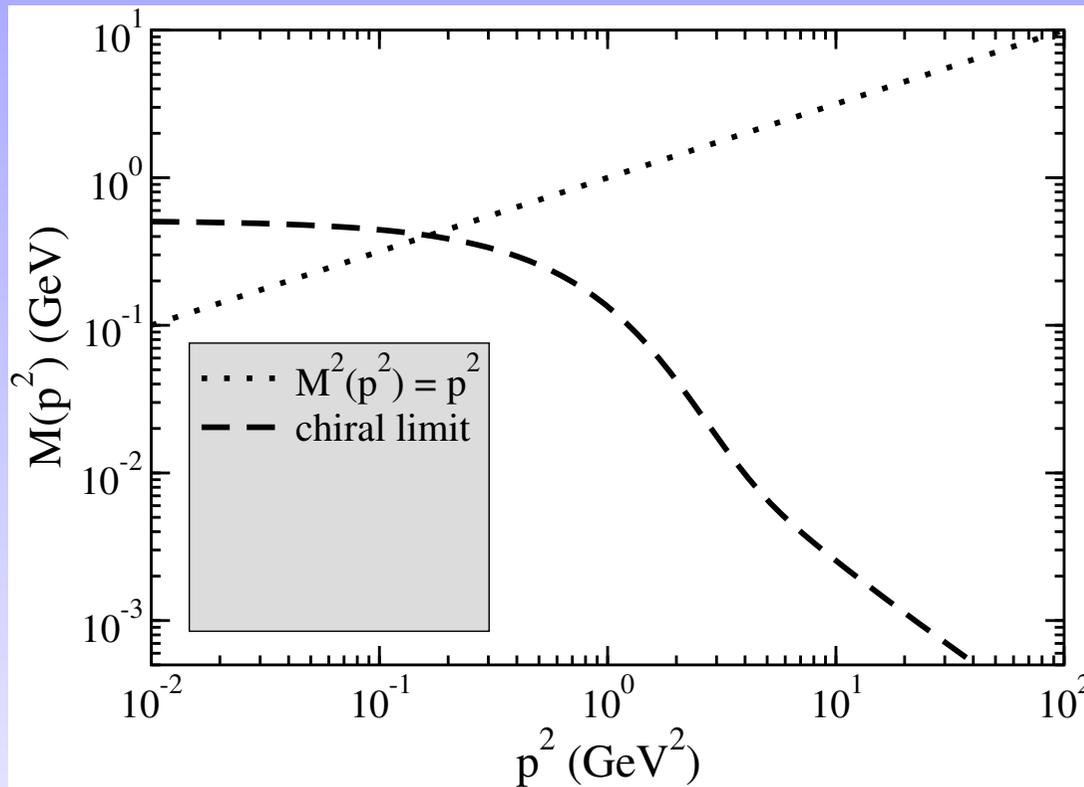
A Challenge



- Infinitely Many Coupled Equations
 - Solutions are Schwinger Functions (Euclidean **Green** Functions)
 - Same VEVs measured in Lattice QCD simulations
- Coupling between equations necessitates truncation
- “Many Body” physics approach
 - Rainbow–Ladder (\sim Hartree–Fock) truncation. Widely explored
 - Systematically improvable.

Nonperturbative Quark Propagator

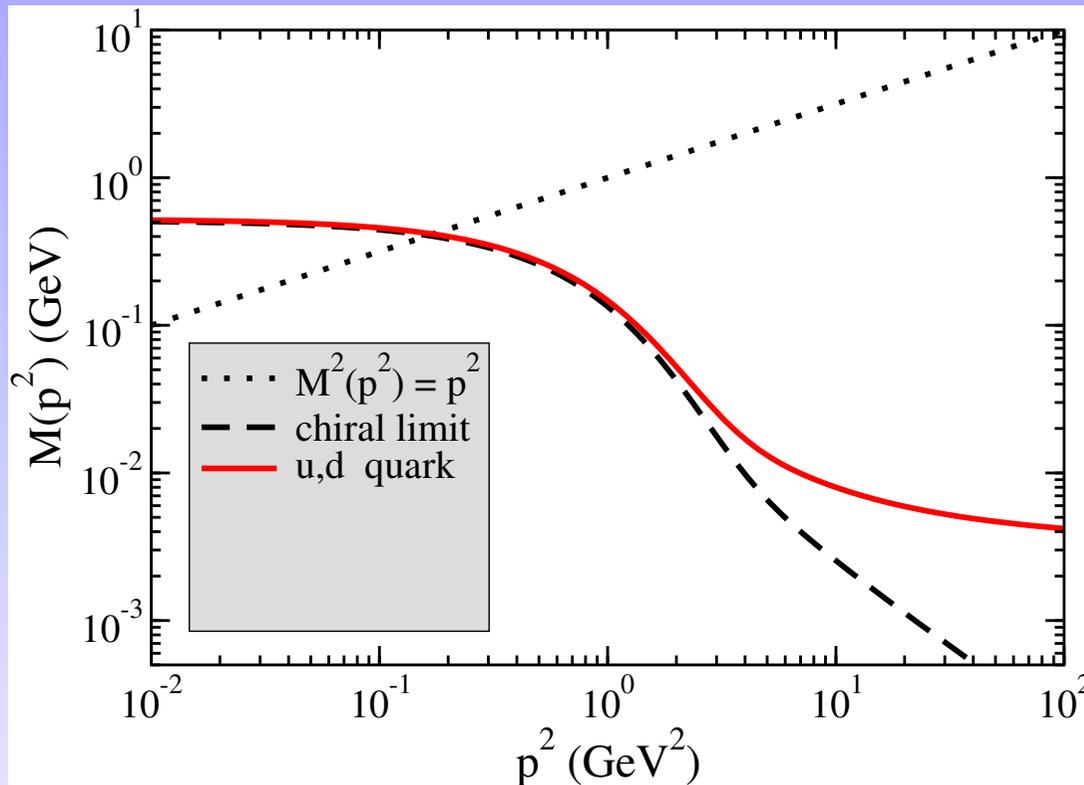
QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$



Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

Nonperturbative Quark Propagator

QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$

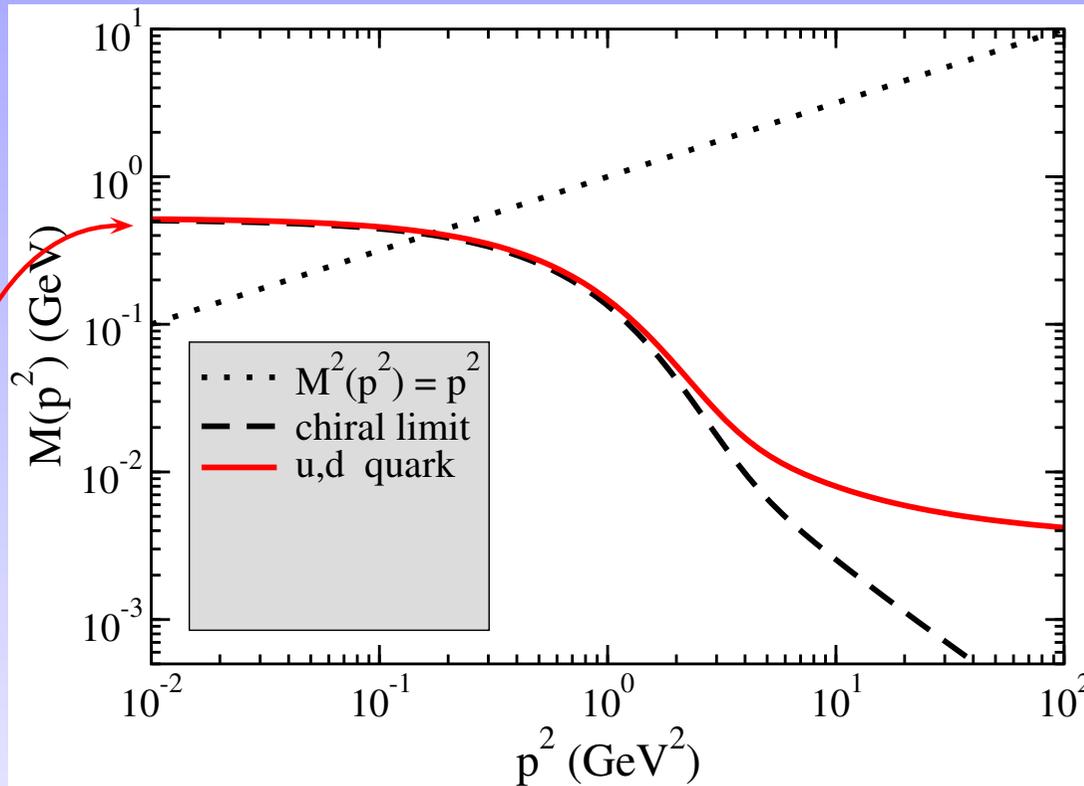


Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

- If $m_q \neq 0$, $M(p^2)$ describes the evolution from a current quark mass at high energies

Nonperturbative Quark Propagator

QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$

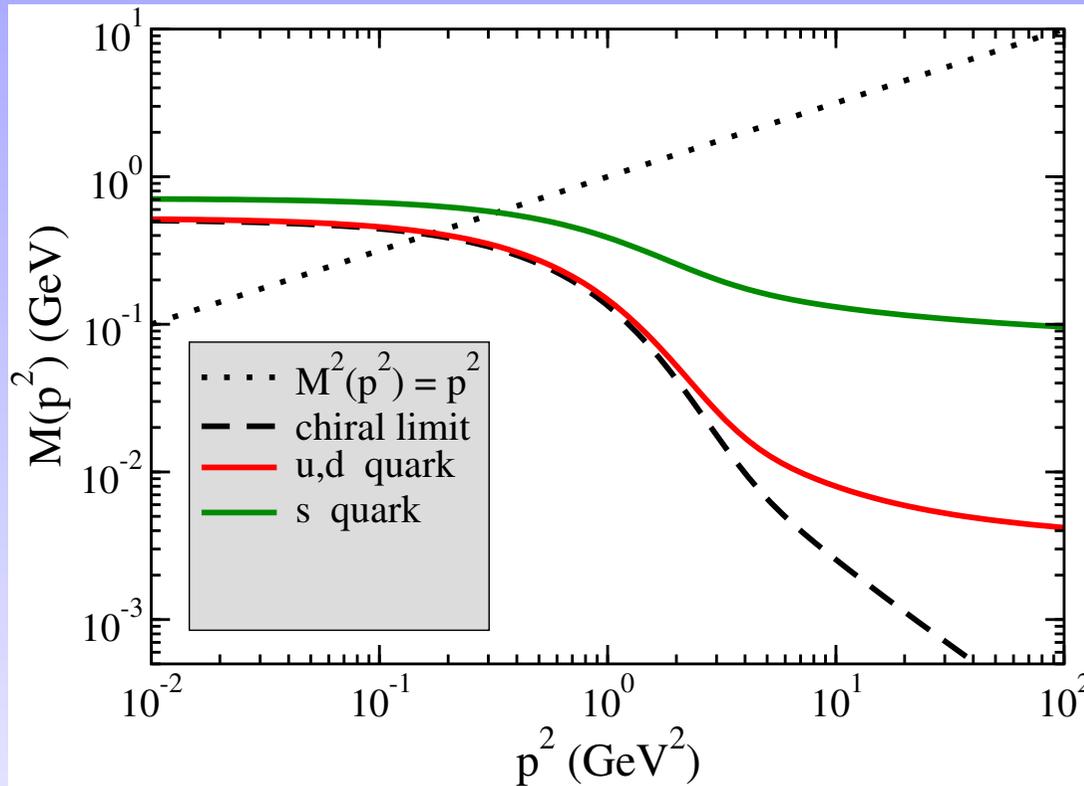


Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

- If $m_q \neq 0$, $M(p^2)$ describes the evolution from a current quark mass at high energies to a constituent-like quark mass at low energies.

Nonperturbative Quark Propagator

QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$



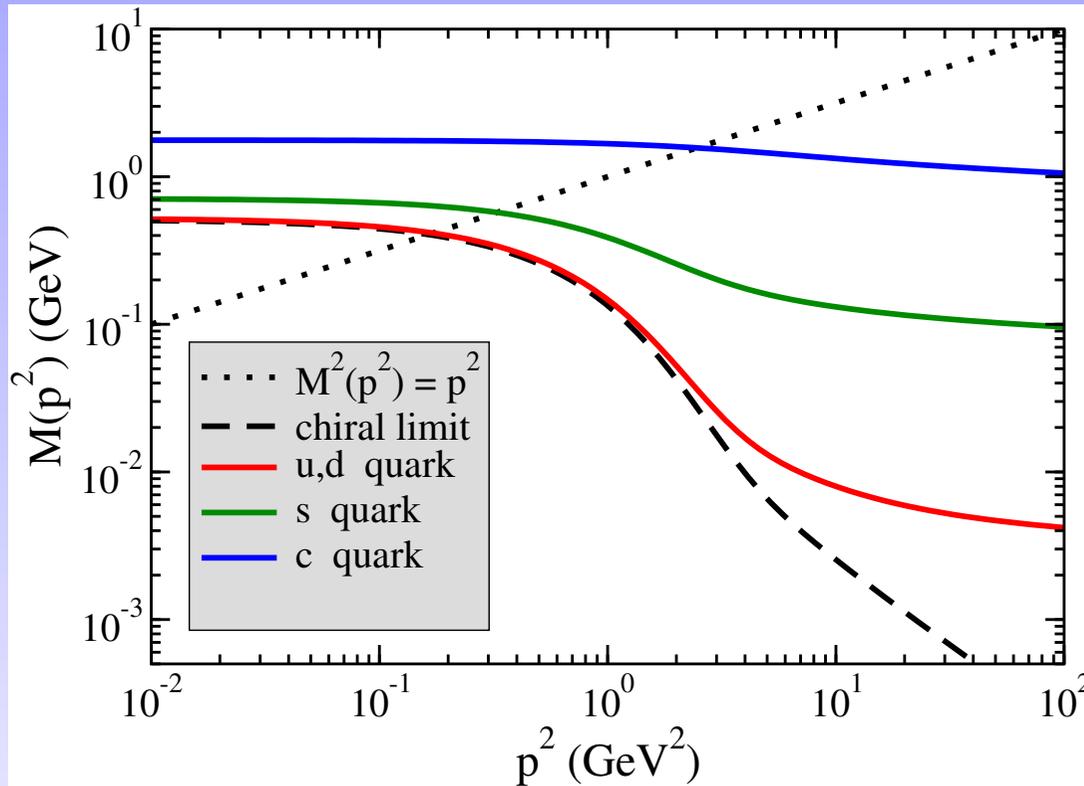
Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

$M(p^2)$ connects nonperturbative QCD with
perturbative QCD



Nonperturbative Quark Propagator

QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$



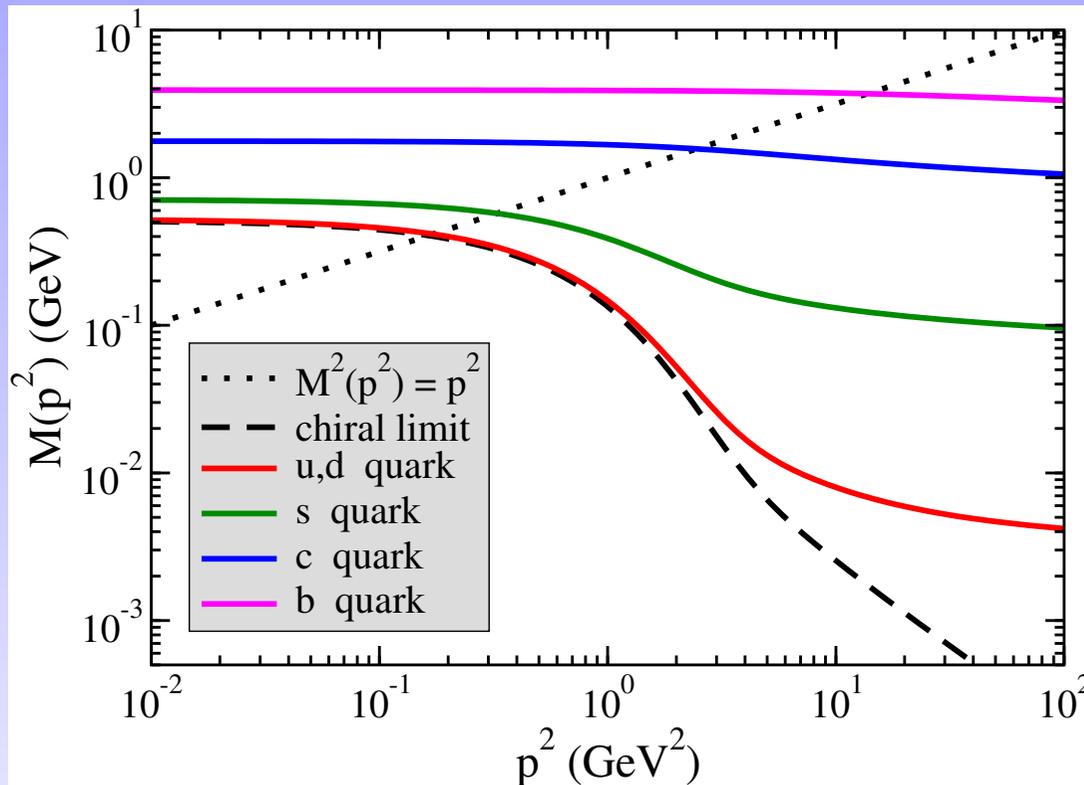
Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

$M(p^2)$ connects nonperturbative QCD with
perturbative QCD



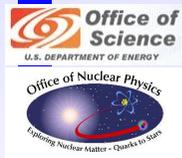
Nonperturbative Quark Propagator

QCD gap equation has solution $M(p^2) \neq 0$ for $m_q = 0$



Adapted from P. Maris &
C.D. Roberts, PRC56,
3369 (1997)
[nucl-th/9708029]

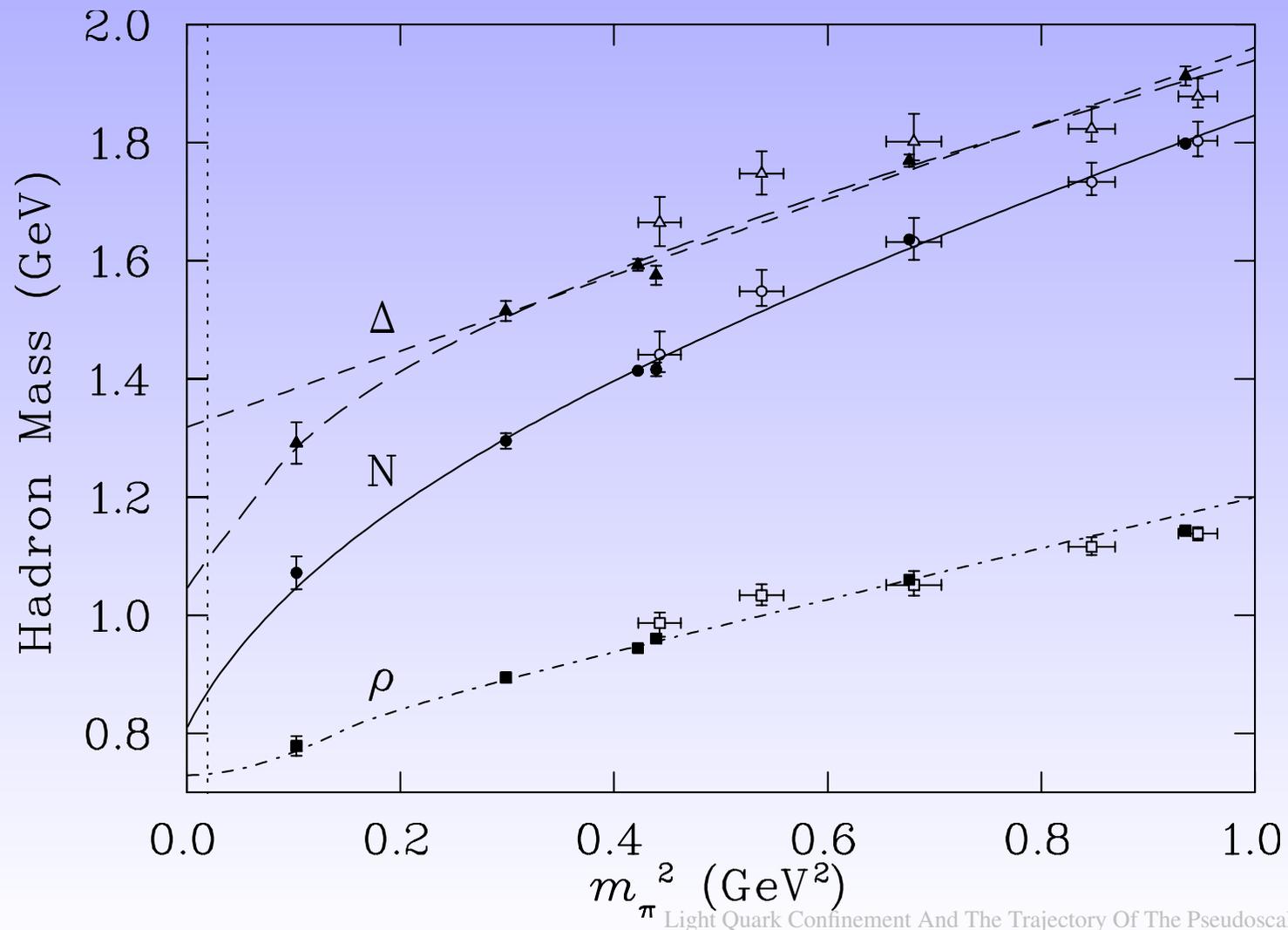
$M(p^2)$ connects nonperturbative QCD with
perturbative QCD



... On the Lattice

Leinweber, Thomas, Tsushima, SVW

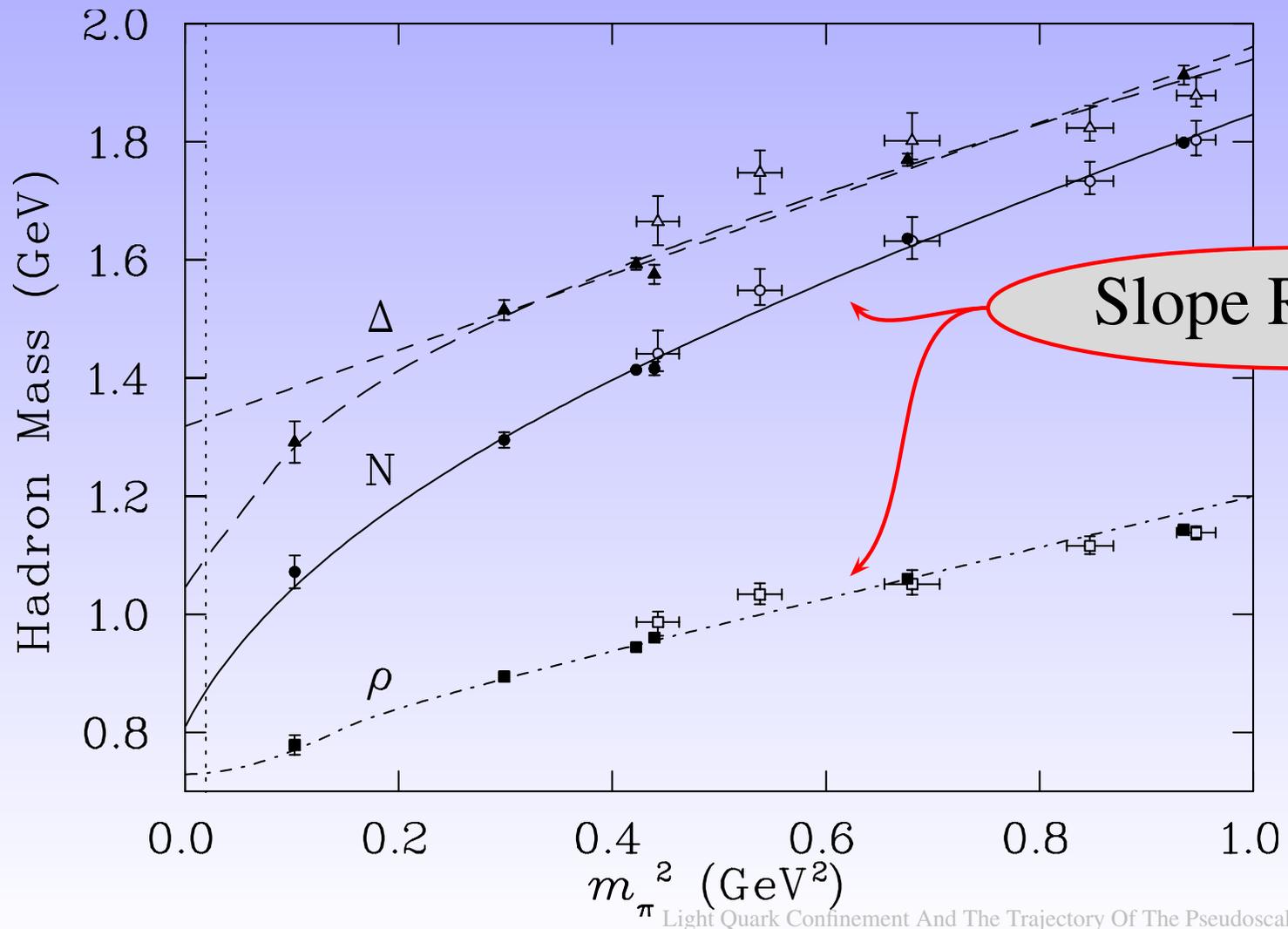
Phys. Rev. **D61**, 074502; Phys. Rev. **D64**, 094502



... On the Lattice

Leinweber, Thomas, Tsushima, SVW

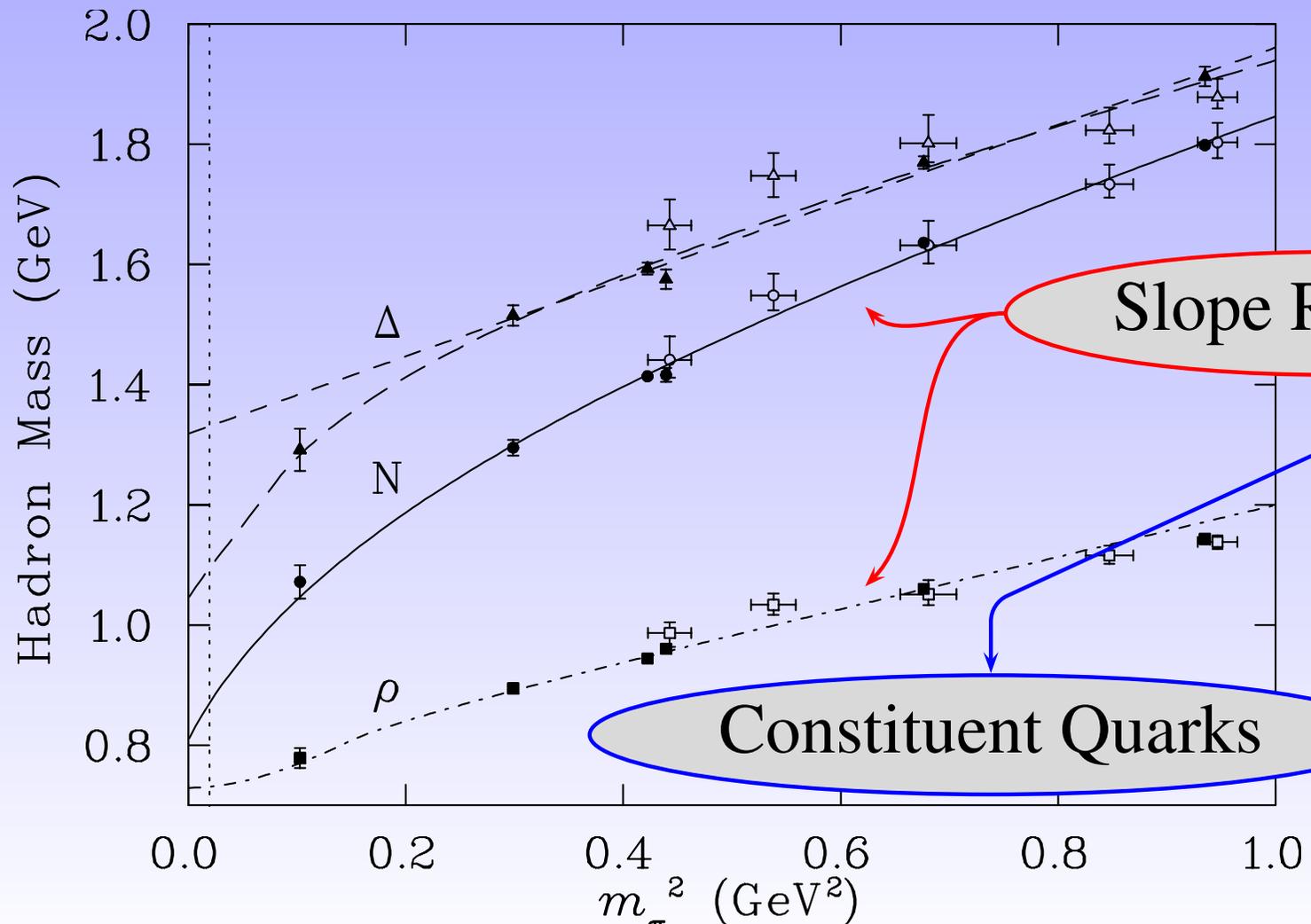
Phys. Rev. **D61**, 074502; Phys. Rev. **D64**, 094502



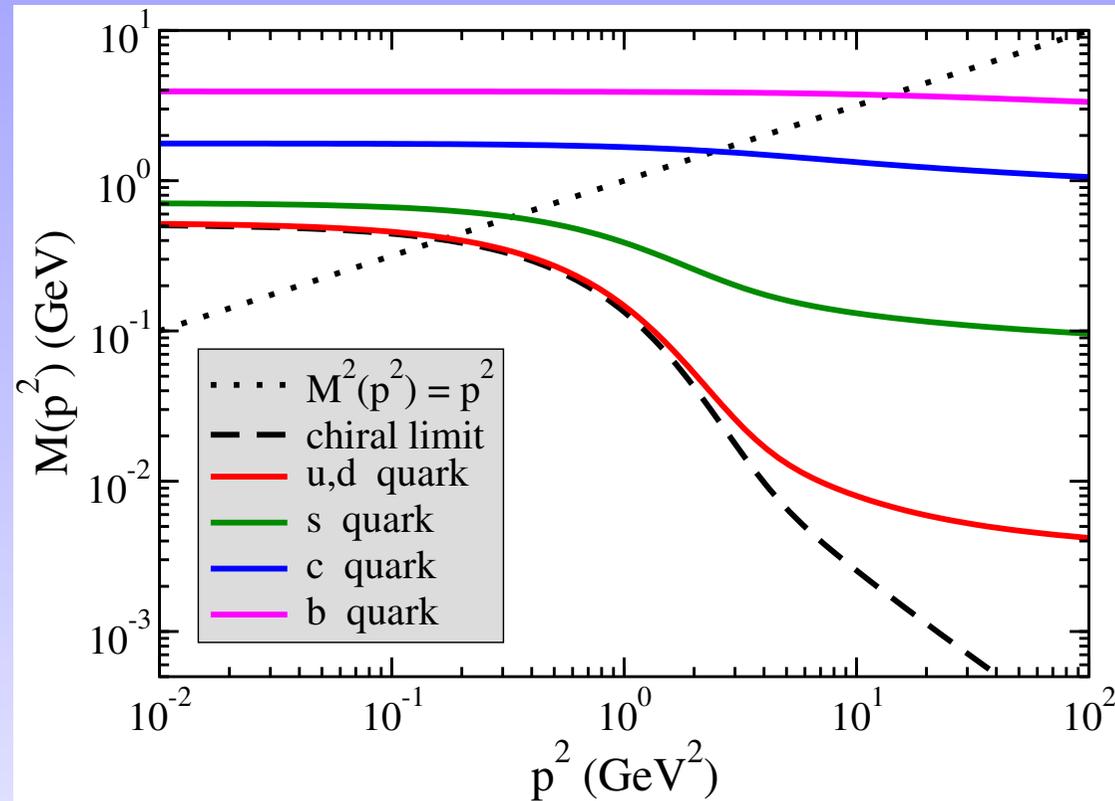
... On the Lattice

Leinweber, Thomas, Tsushima, SVW

Phys. Rev. **D61**, 074502; Phys. Rev. **D64**, 094502



Nonperturbative Mass Function

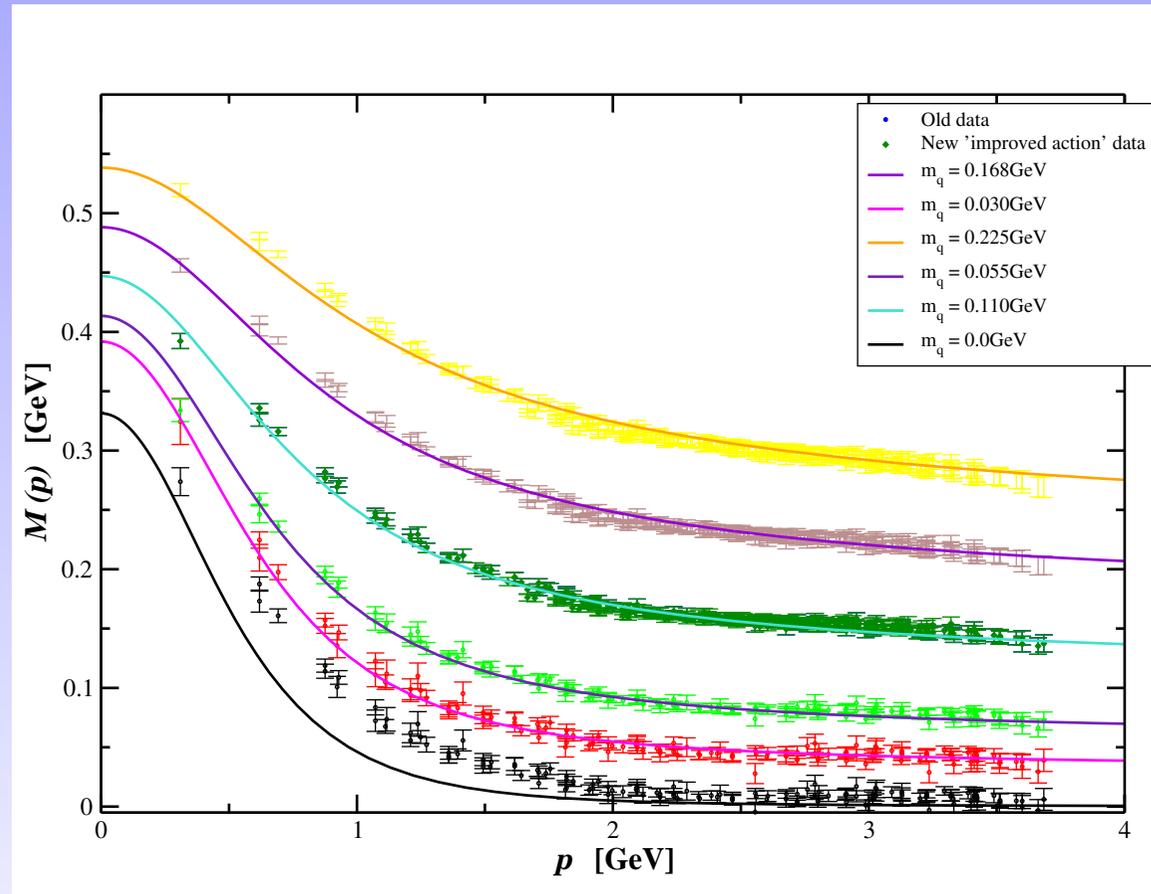


Chiral Symmetry and its dynamical breaking
is important property for light **Quarks**

Comparison with the Lattice

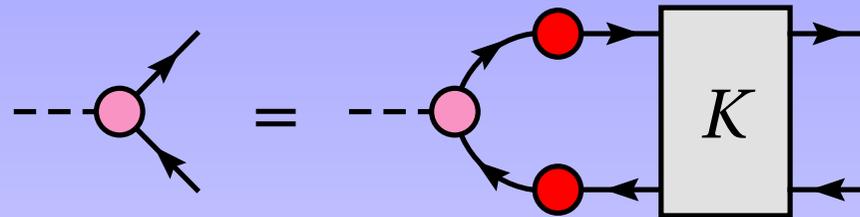
- One parameter model provides persuasive results.

DSE: Bhagwat,
Pichowsky,
Roberts, Tandy
nucl-th/0304003
Lattice:
Bowman, Heller,
Leinweber, Williams
hep-lat/0209129



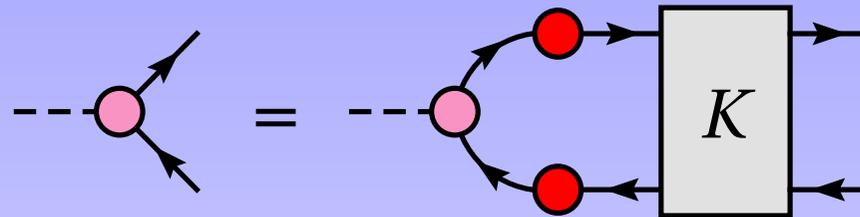
Hadrons: Bound States (BSE)

Bethe–Salpeter Equation



Hadrons: Bound States (BSE)

Bethe–Salpeter Equation



- Axial-Vector Ward–Takahashi identity relates kernel of BSE and DSE

My interest



- Coming from a (pseudo-) Lattice background

My interest



- Coming from a (pseudo-) Lattice background
 - I *like* observables

My interest



- Coming from a (pseudo-) Lattice background
 - I *like* observables
- Ground state mesons thoroughly investigated
 - ⇒ Excited States

Why?

SPIRES: Almost 70% of papers with “*excited meson(s)*” in the title published in the last 10 years.



Collaborators



ARGONNE: Craig Roberts, Arne Höll

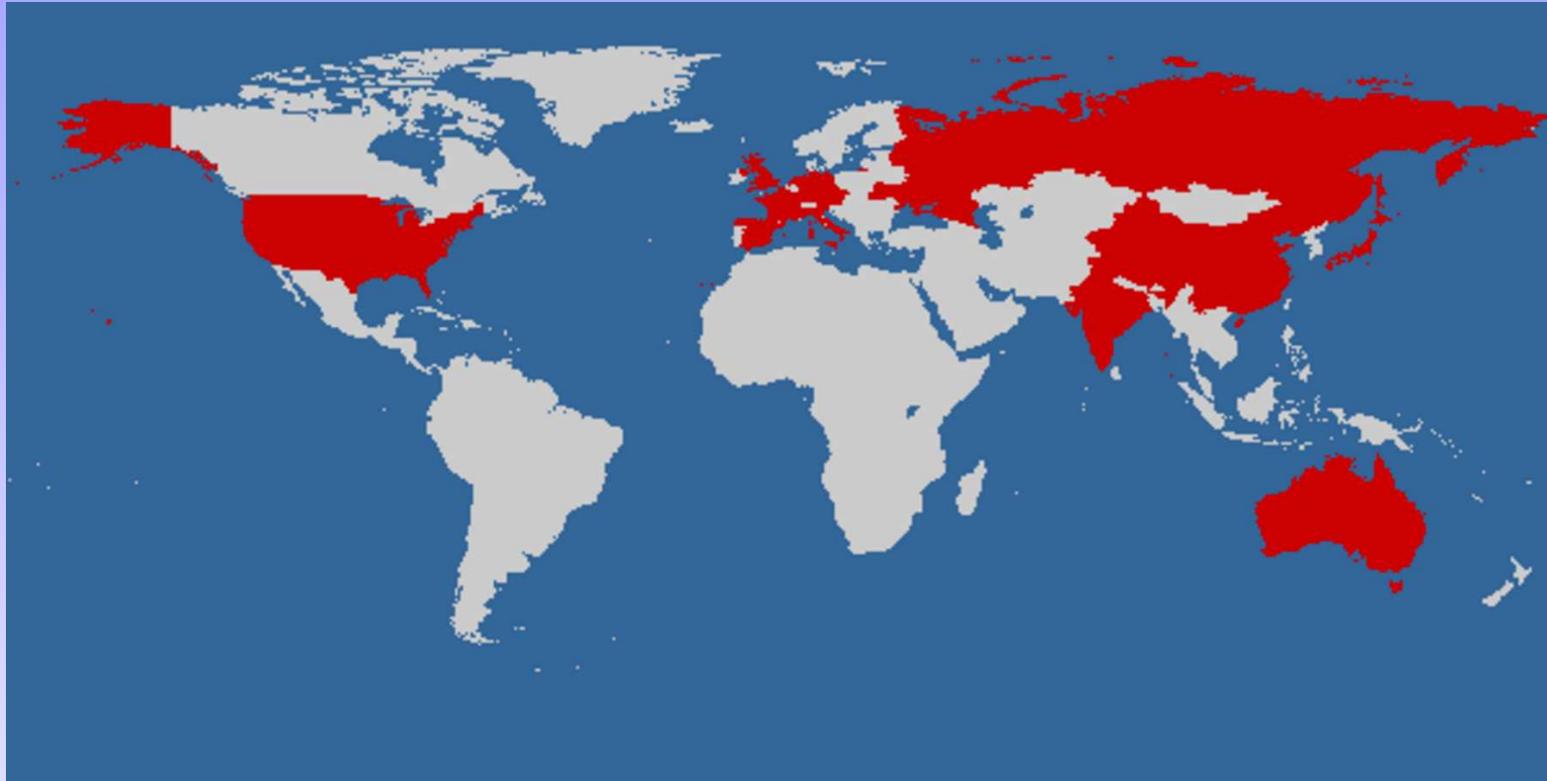
PITTSBURG UNIV: Pieter Maris

GRAZ UNIV: Andreas Krassnigg

KENT STATE UNIV: Peter Tandy
(Mandar Bhagwat)



Recent DSE Publications



Last 25 papers listed on SPIRES with “*DSE*” in the title...



This Study: Pions

- Solve inhomogeneous vertex DSE.



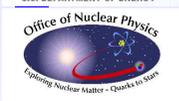
This Study: Pions

- Solve inhomogeneous vertex DSE.
- Numerical solution exists for both TIMELIKE and SPACELIKE momentum.



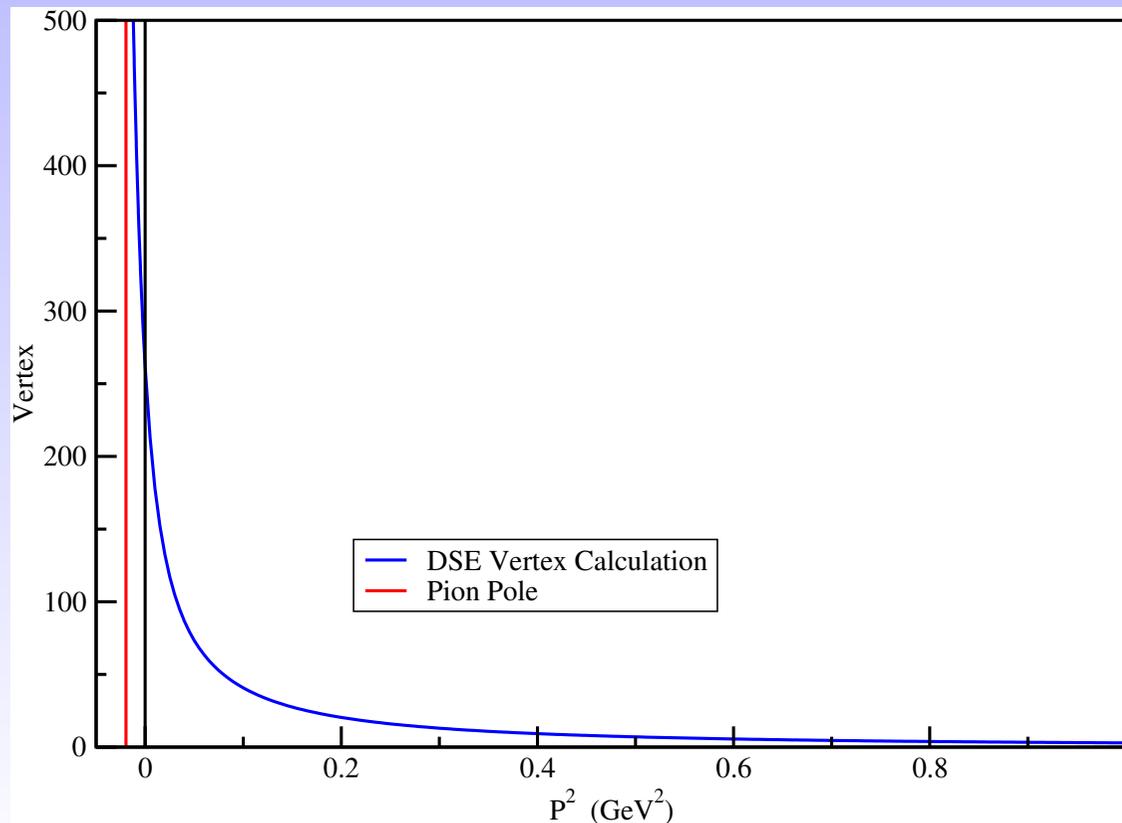
This Study: Pions

- Solve inhomogeneous vertex DSE.
- Numerical solution exists for both TIMELIKE and SPACELIKE momentum.
- Poles in timelike region \rightarrow bound states

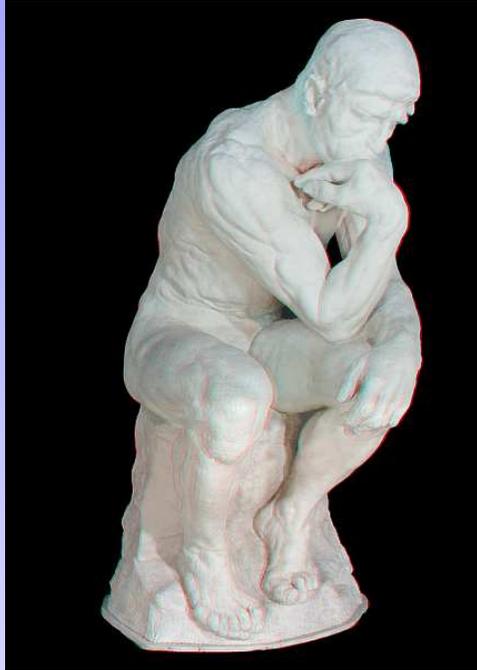


This Study: Pions

- Solve inhomogeneous vertex DSE.
- Numerical solution exists for both TIMELIKE and SPACELIKE momentum.
- Poles in timelike region \rightarrow bound states



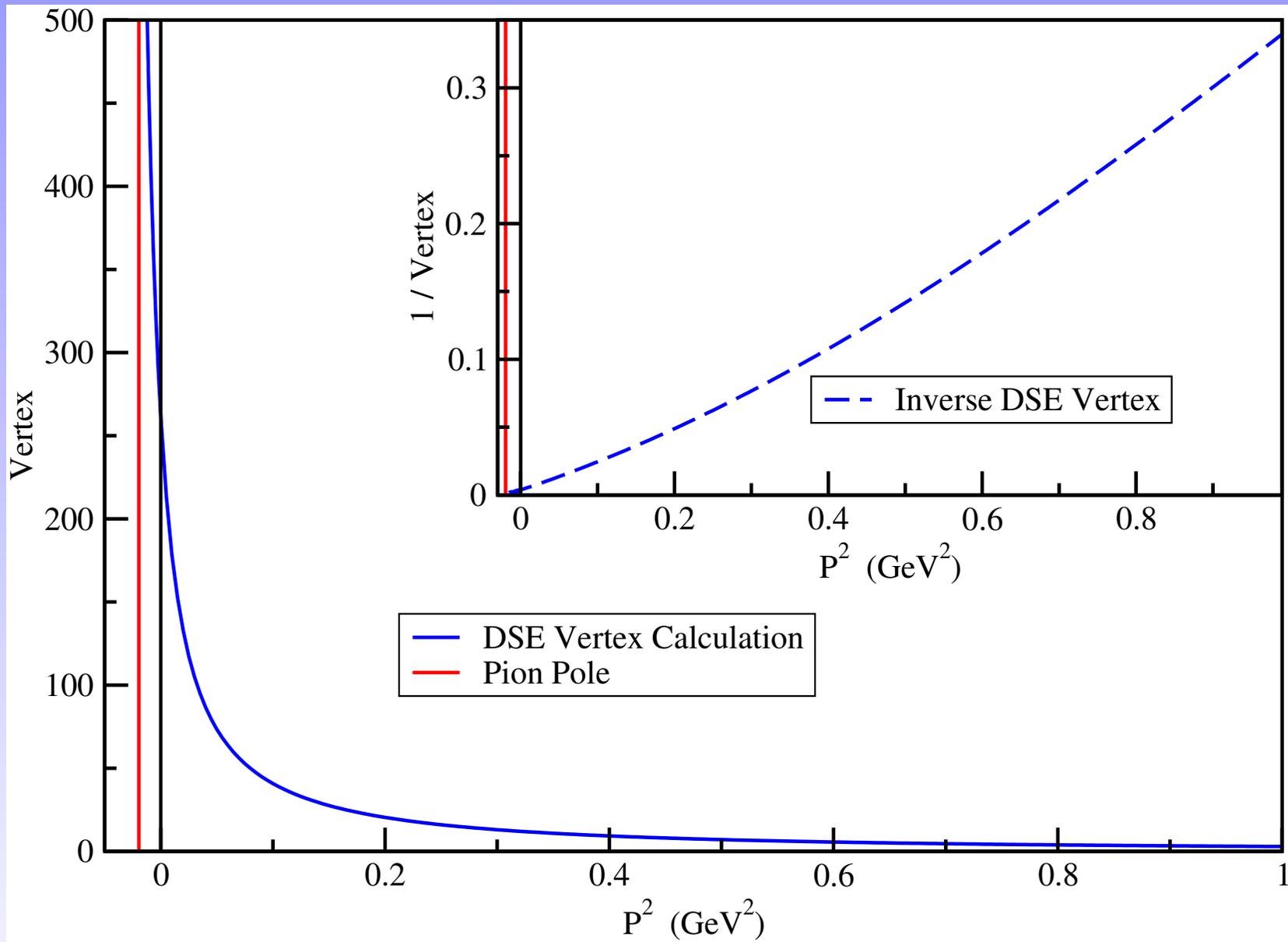
Numerical Reality

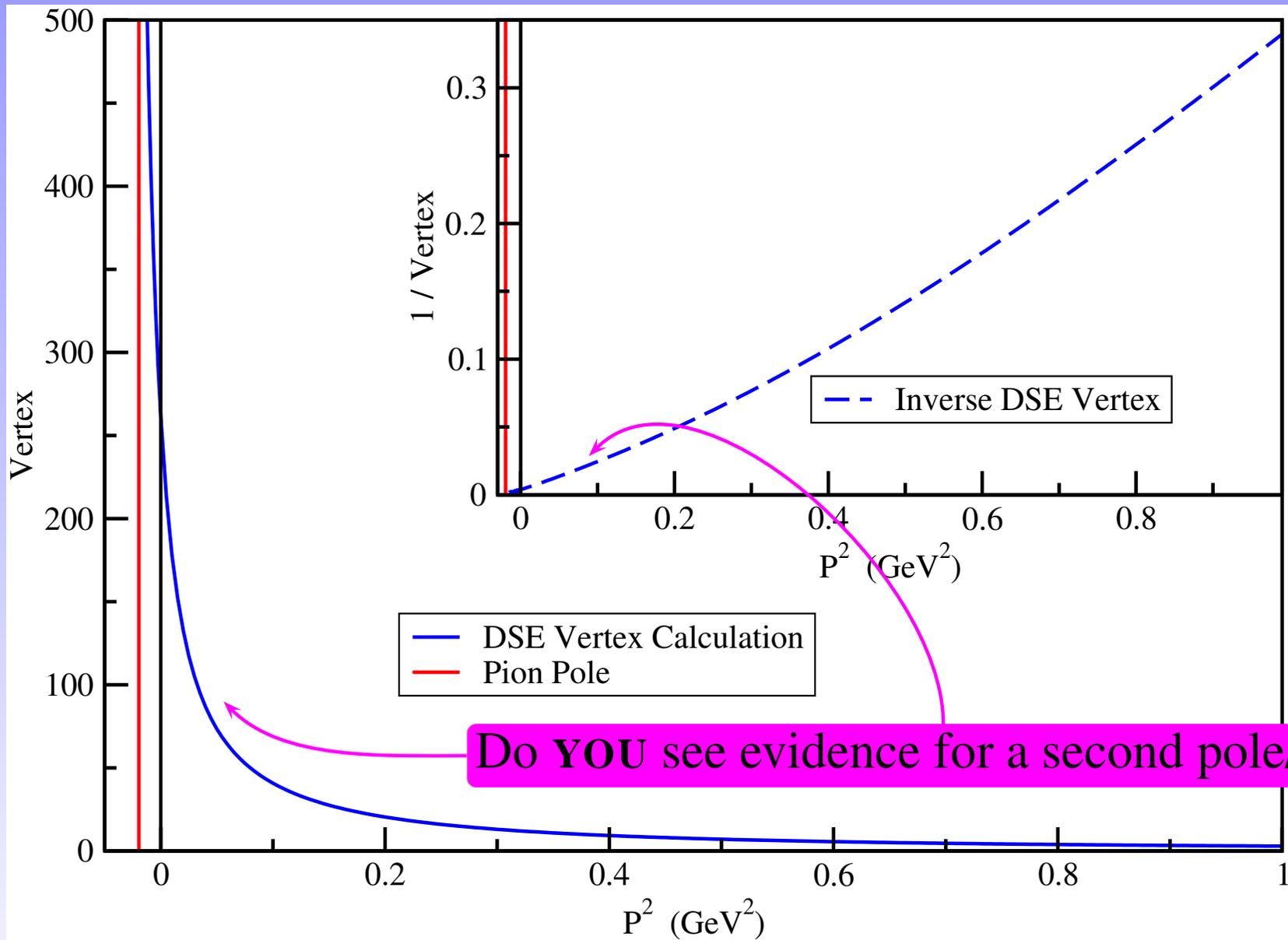


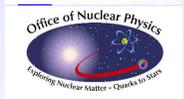
☹ Finding exact pole location
is HARD

☺ Finding zeros is EASY

So...invert the data!







What do we know about calculating excited states?

What do we *think* we know?



How many excited states can one expect to find when you have some vertex correlator?

Simple Model Investigation

Naïvely the spectral model of a 3-point function is a sum of bound states.

Make a simple model for a vertex:

- A sum of monopoles

$$V(\vec{P}) = a + \sum_i \frac{c_i}{\vec{P}^2 + m_i^2}$$

- m_i : Mass of bound state i
- c_i : Residue of bound state i

Use some physics

What do we know about m_i and c_i ?

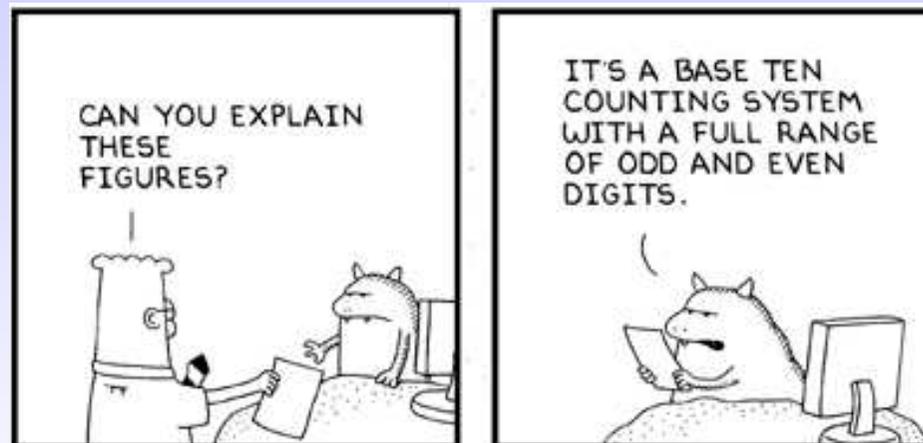


Use some physics

What do we know about m_i and c_i ?

Mass

- Masses of the bound states
- PDG publishes these for the REAL WORLD



Use some physics

What do we know about m_i and c_i ?

Mass

- Masses of the bound states
- PDG publishes these for the REAL WORLD

Residue

- Related to the decay constant of the bound state, i.e. f_π
- We **know** that these alternate in sign

Höll, Krassnigg, Roberts nucl-th/0406030



Visualise the problem

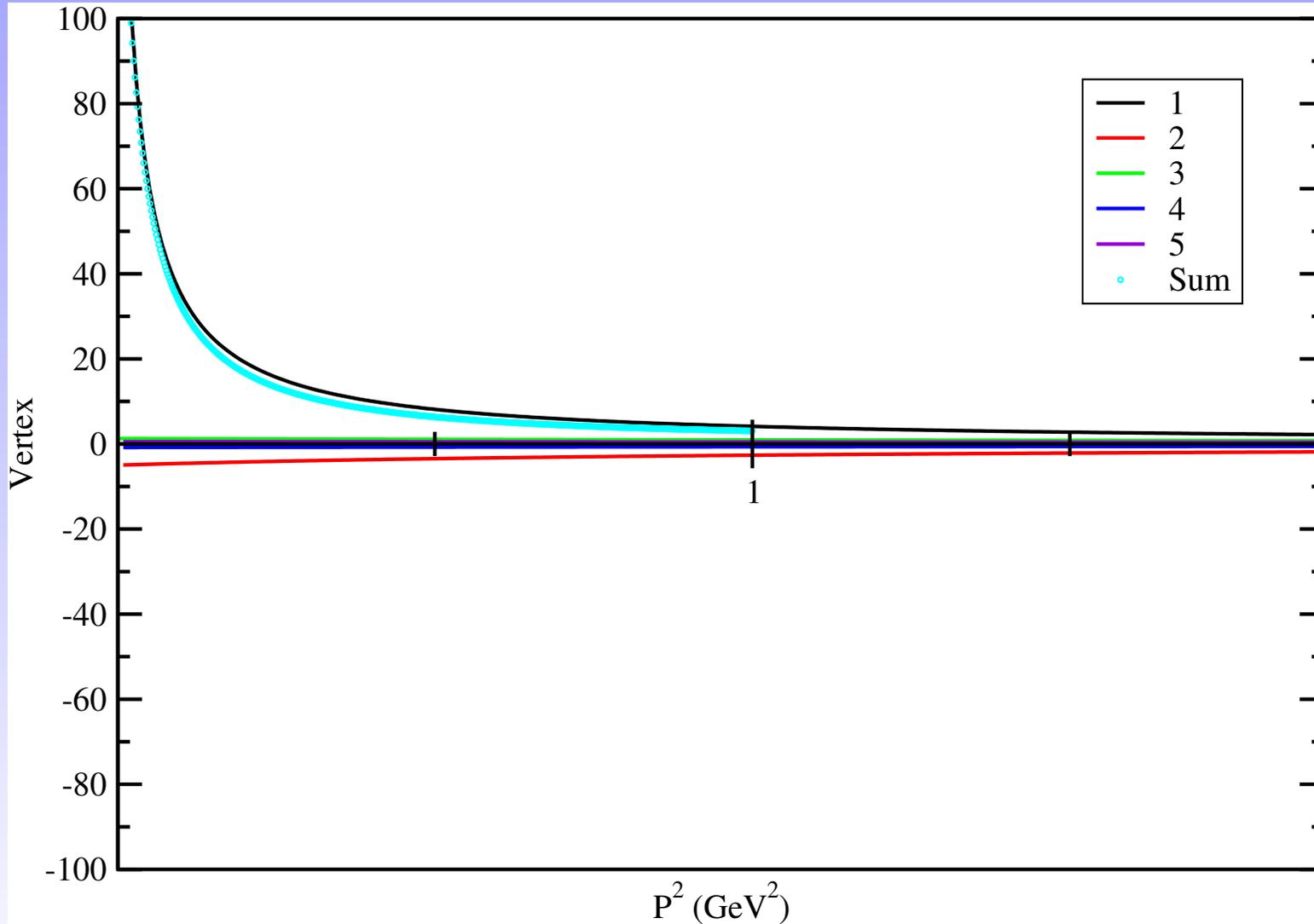
Mass	Residue
0.14	4.23
1.06	-5.6
1.72	3.82
2.05	-3.45
2.2	2.8

Masses are motivated by the PDG.

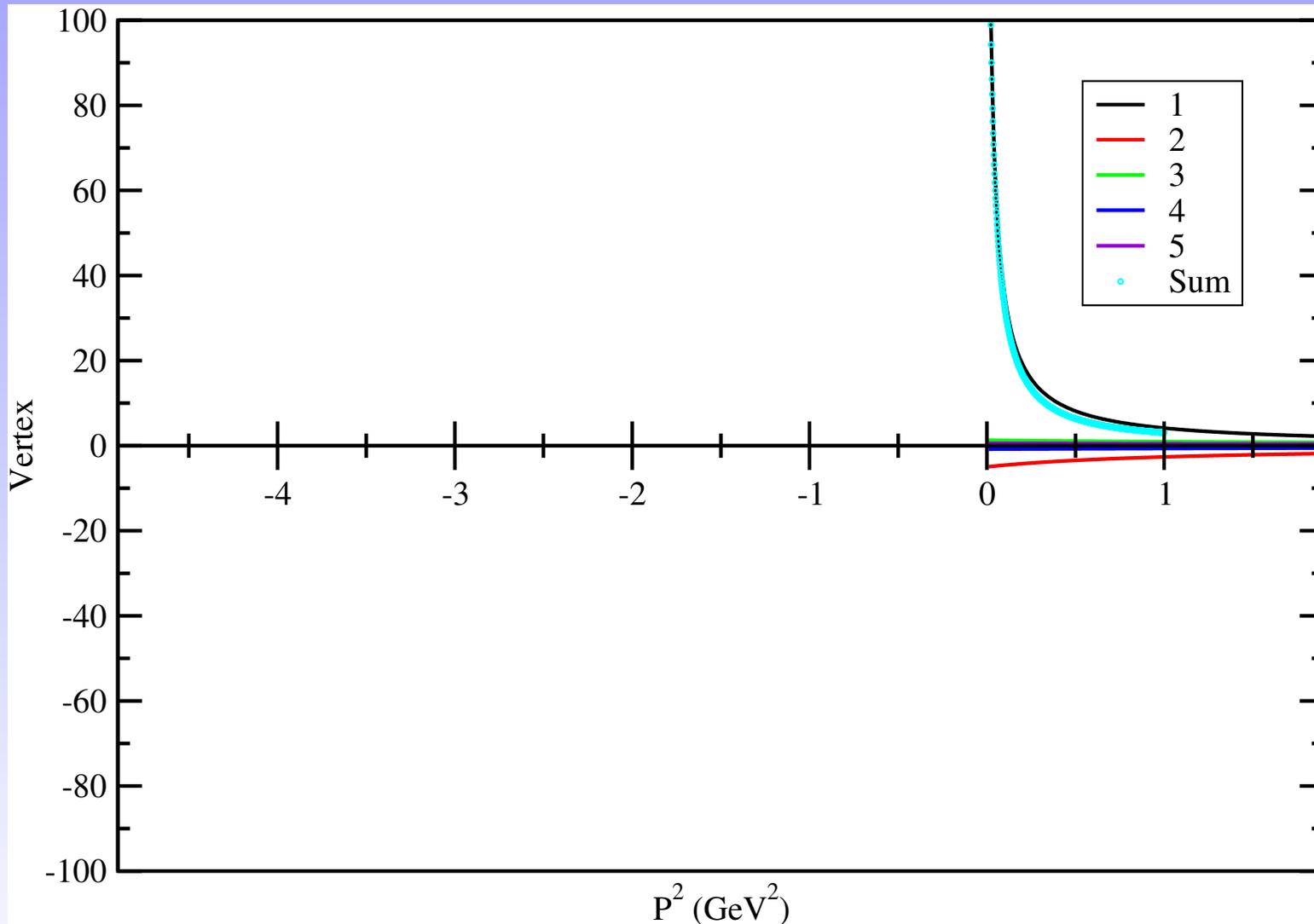
Residues are all of the same magnitude.



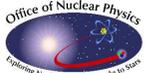
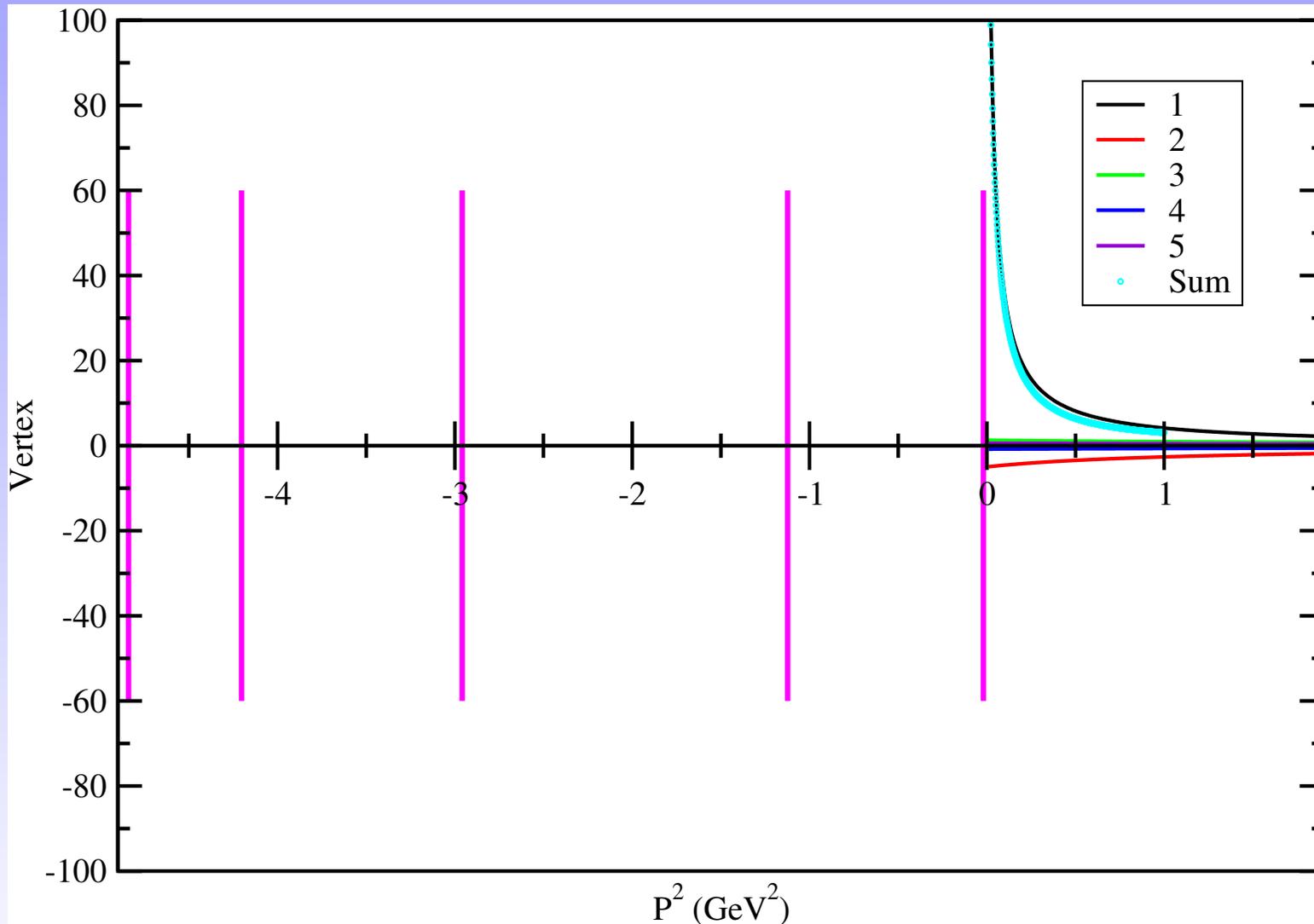
The Simple Model



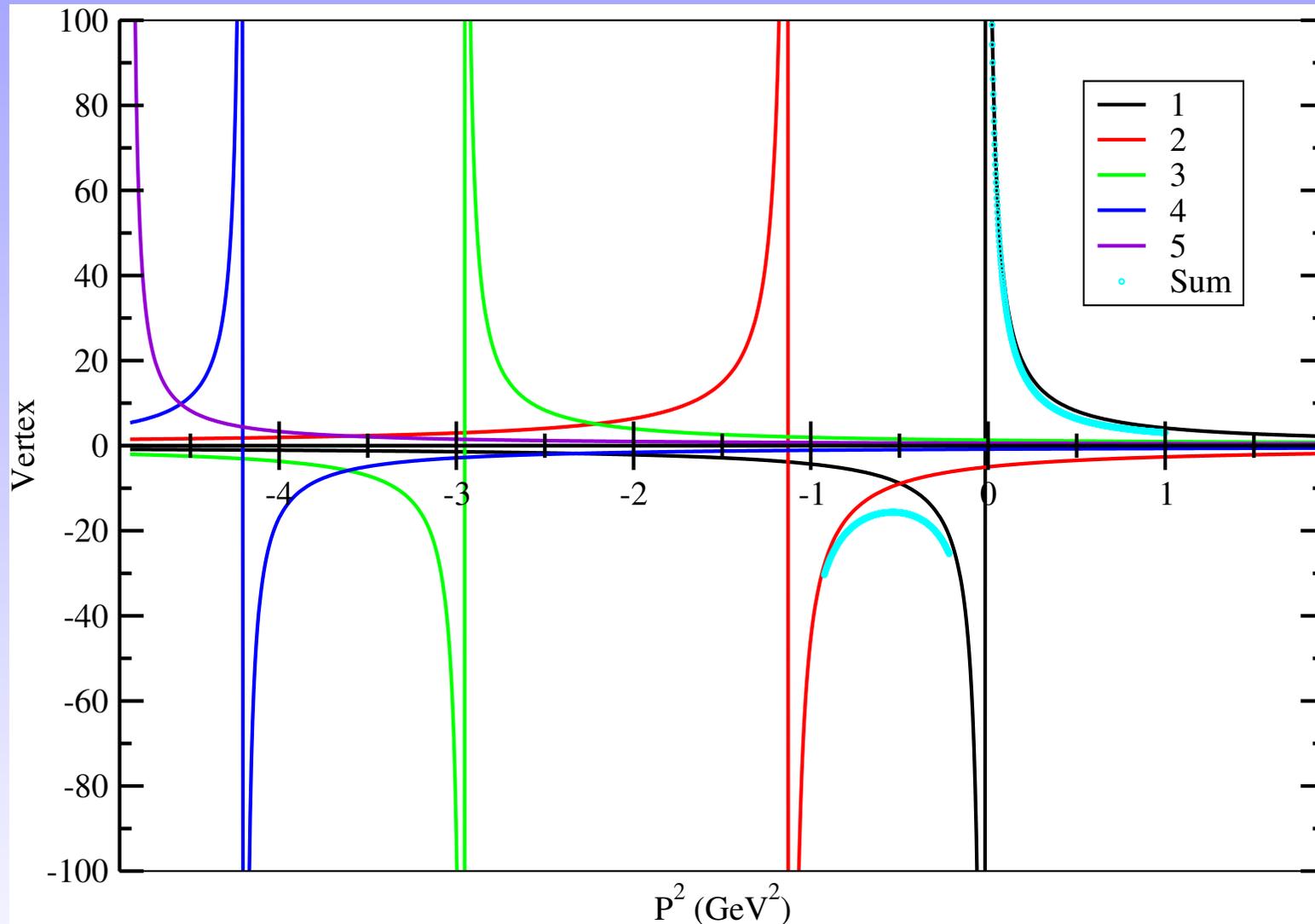
The Simple Model



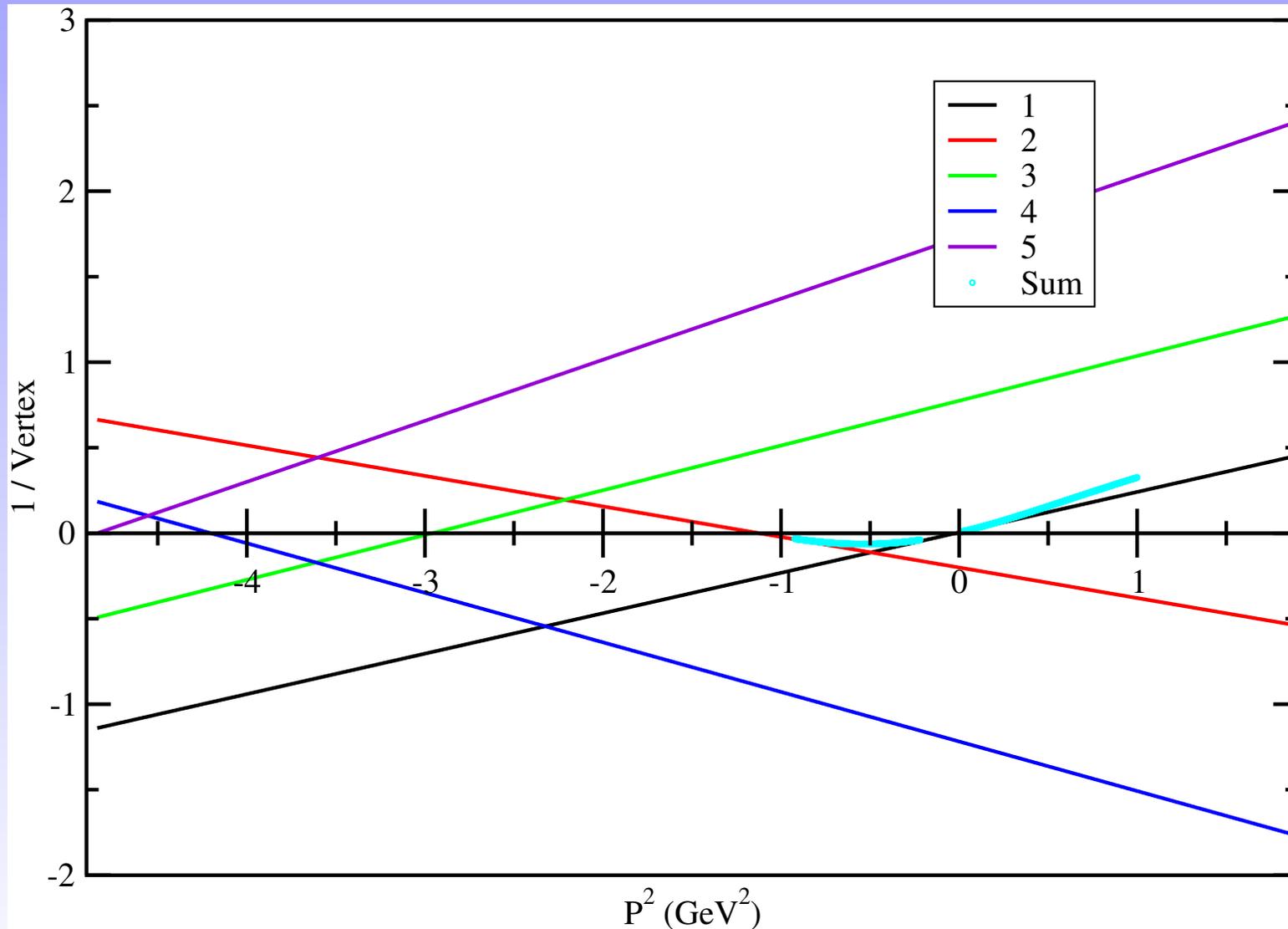
The Simple Model



The Simple Model



Invert for zeros



Fitting the Data

Assume we know nothing about the source of the data...



Fitting the Data

Assume we know nothing about the source of the data...

Fit a Padé to data:

- We know the data has zeros
- Little else is known about the functional form i.e. what does the background look like?

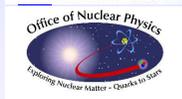
So choose a generic form:

$$f(P^2) = \frac{a_0 + a_1 P^2 + a_2 P^4 + \dots + a_n P^{2n}}{1 + a_{n+1} P^2 + a_{n+2} P^4 + \dots + a_{2n} P^{2n}}$$

Interpreting the solution

The “data” has an unknown number of bound states contributing.

How reliable are the n -zeros we find?

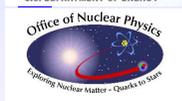


Interpreting the solution

The “data” has an unknown number of bound states contributing.

How reliable are the n -zeros we find?

- The contribution to the vertex from the bound states falls as $1/m^2$
⇒ Clear hierarchy

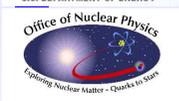


Interpreting the solution

The “data” has an unknown number of bound states contributing.

How reliable are the n -zeros we find?

- The contribution to the vertex from the bound states falls as $1/m^2$
⇒ Clear hierarchy
- At **BEST** the first $n - 1$ solutions would be reliable



Interpreting the solution

The “data” has an unknown number of bound states contributing.

How reliable are the n -zeros we find?

- The contribution to the vertex from the bound states falls as $1/m^2$
⇒ Clear hierarchy
- At **BEST** the first $n - 1$ solutions would be reliable
- The last solution would contain all the remaining physics

How?

The fitting function:

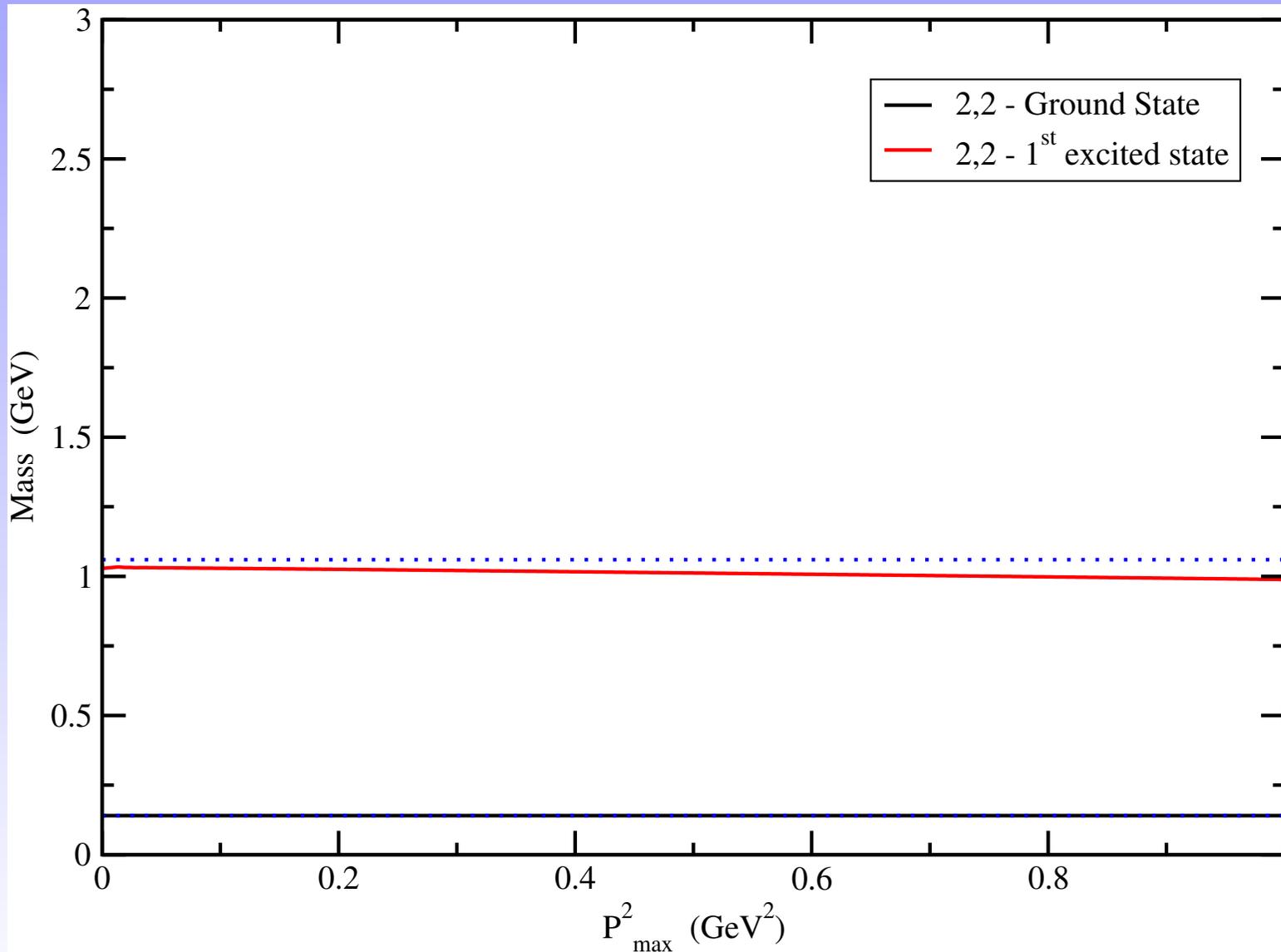
$$f(P^2) = \frac{a_0 + a_1 P^2 + a_2 P^4 + \dots + a_n P^{2n}}{1 + a_{n+1} P^2 + a_{n+2} P^4 + \dots + a_{2n} P^{2n}}$$

Have data for $P^2 \in (0, 1] \text{ GeV}^2$.

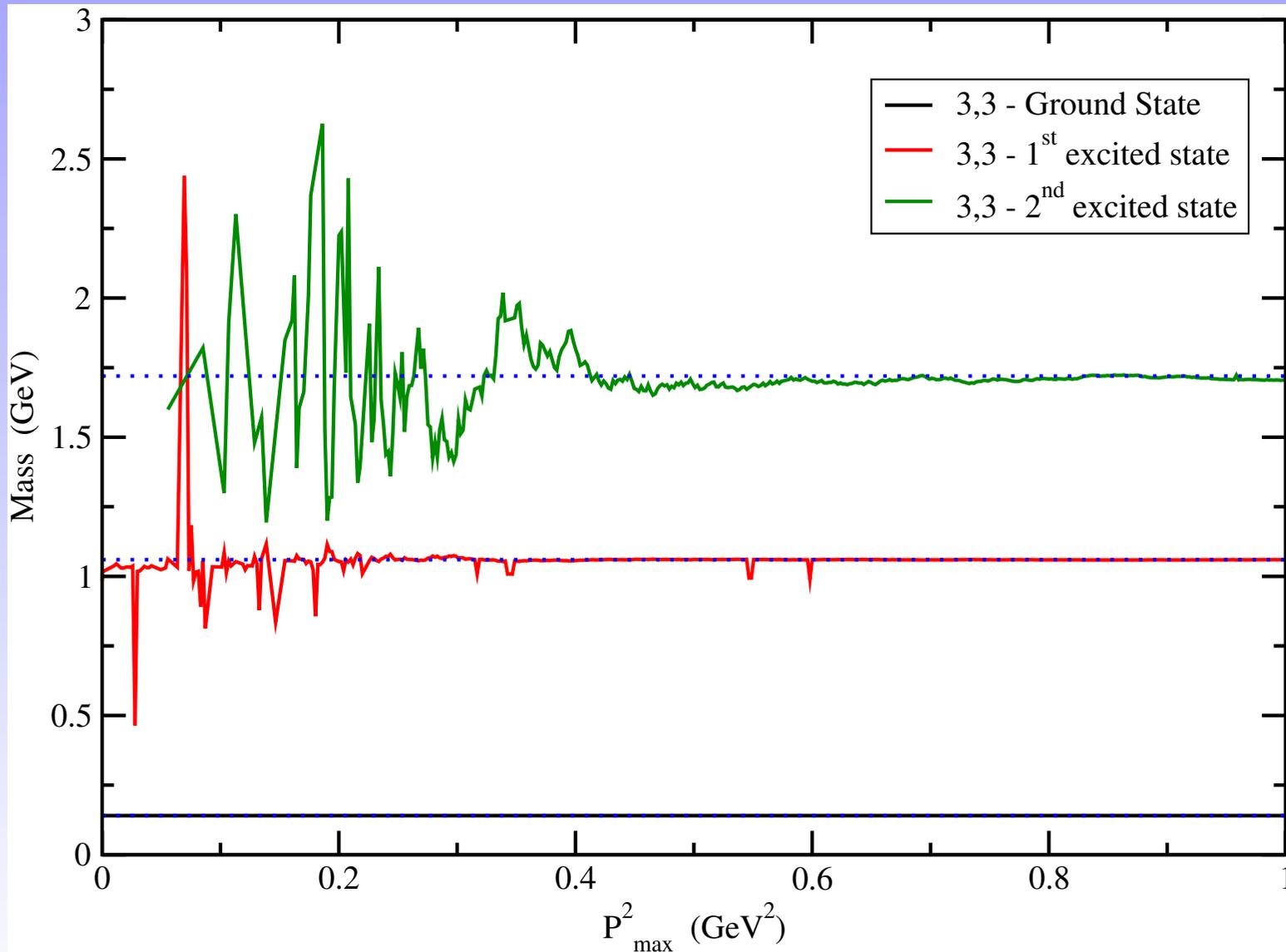
1. Fit data in range $P^2 = (0, P_{max}^2]$.
2. Increase P_{max}^2 and repeat fit.

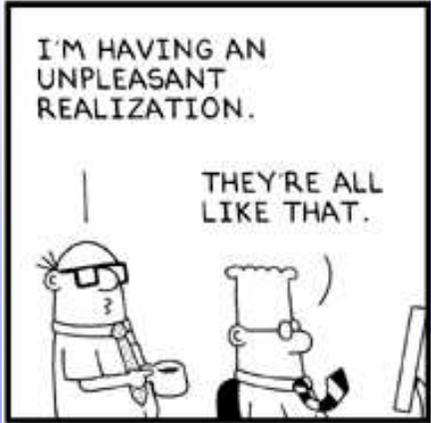


Results — Masses

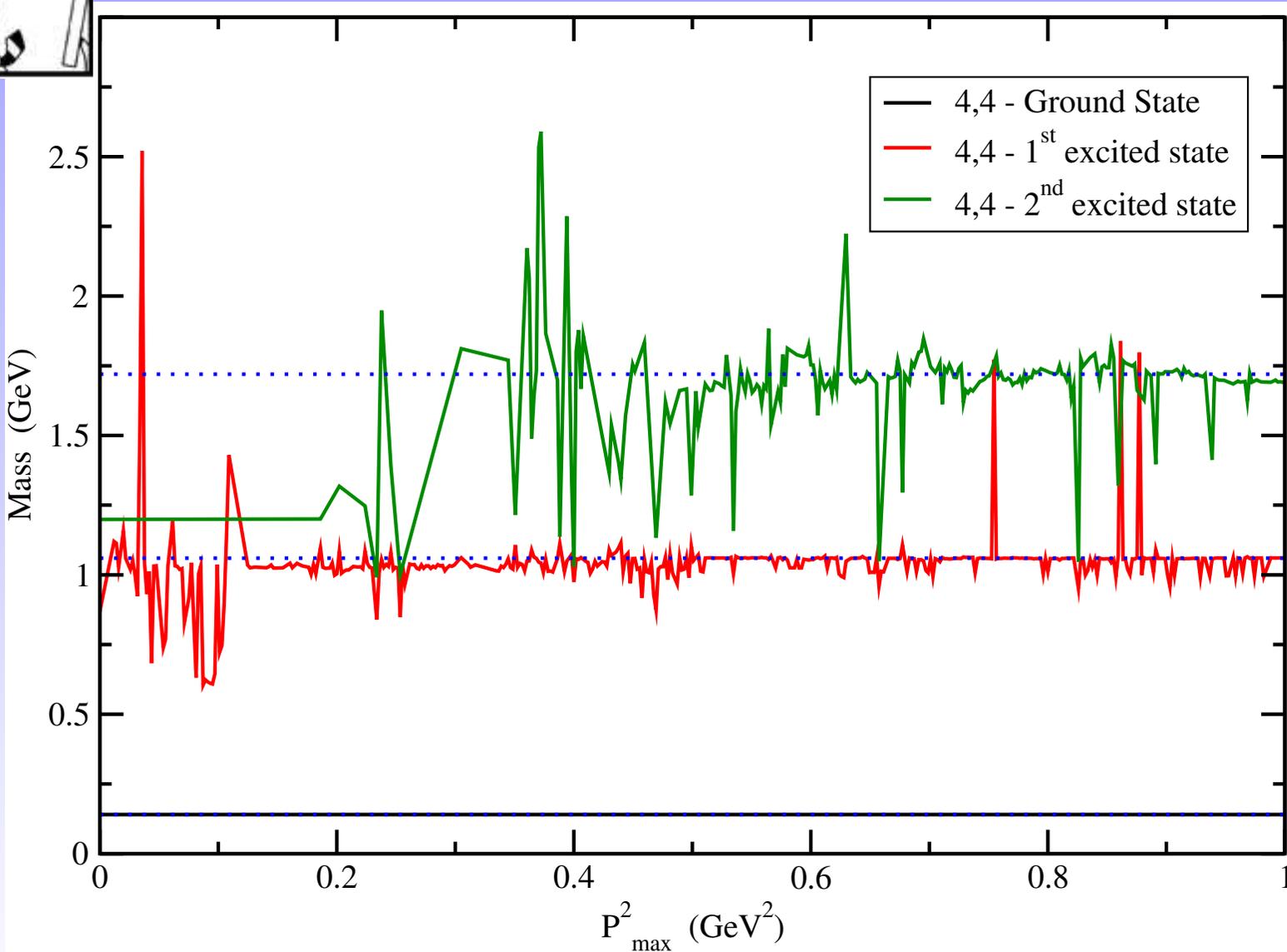


Results — Masses

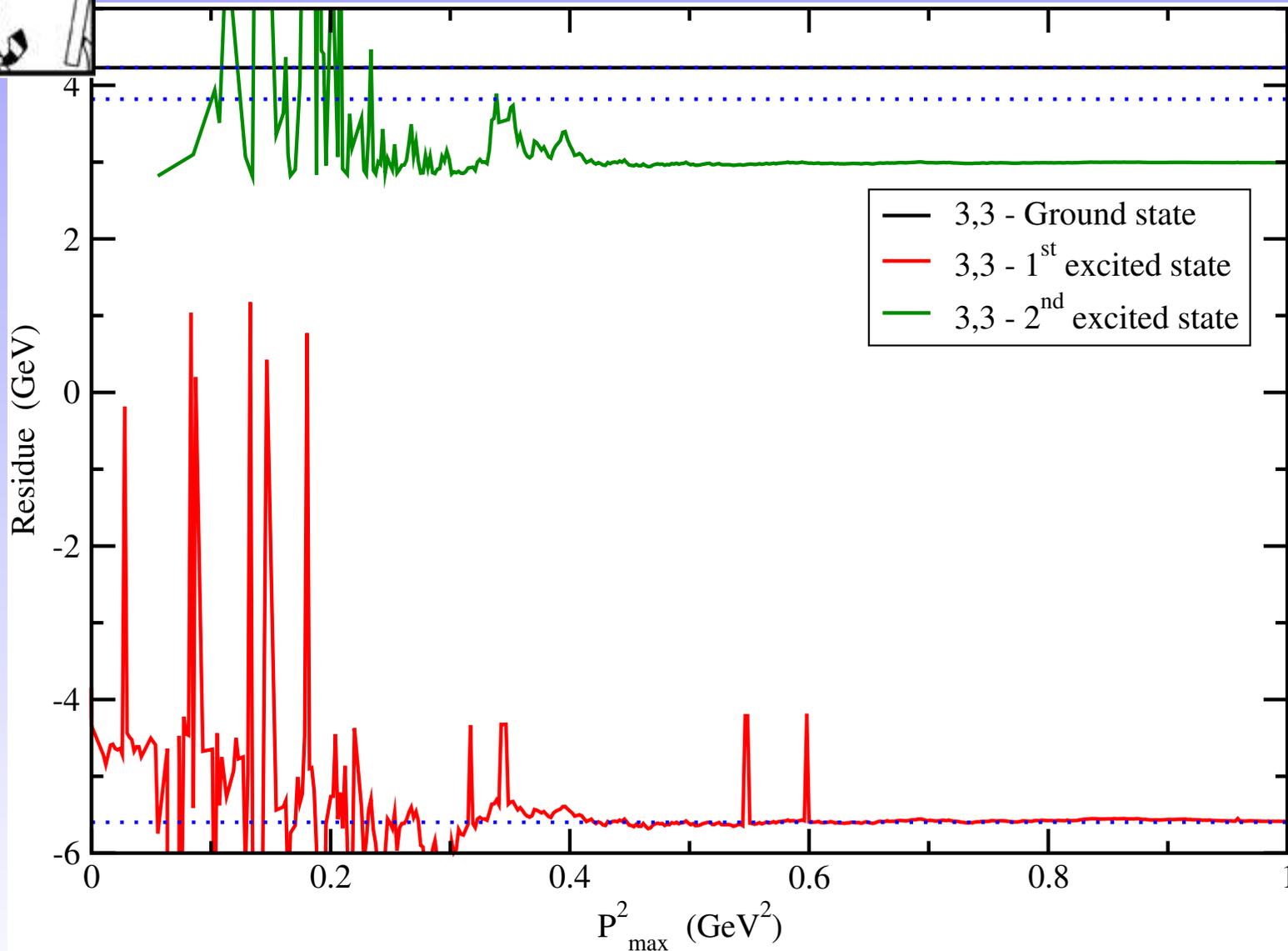
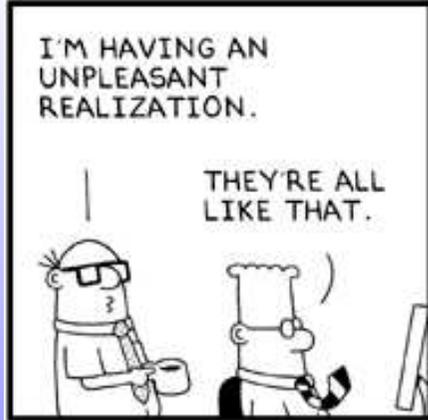




Results — Masses



Results — Residues



The Glass is Half Full

Optimistically there is a chance to get some information about excited states with only spacelike data...

	Mass	Residue
Ground State:		
First Excited State:		
Second Excited State:		

Lattice QCD

A similar investigation in the Lattice QCD mindset is relatively easy:

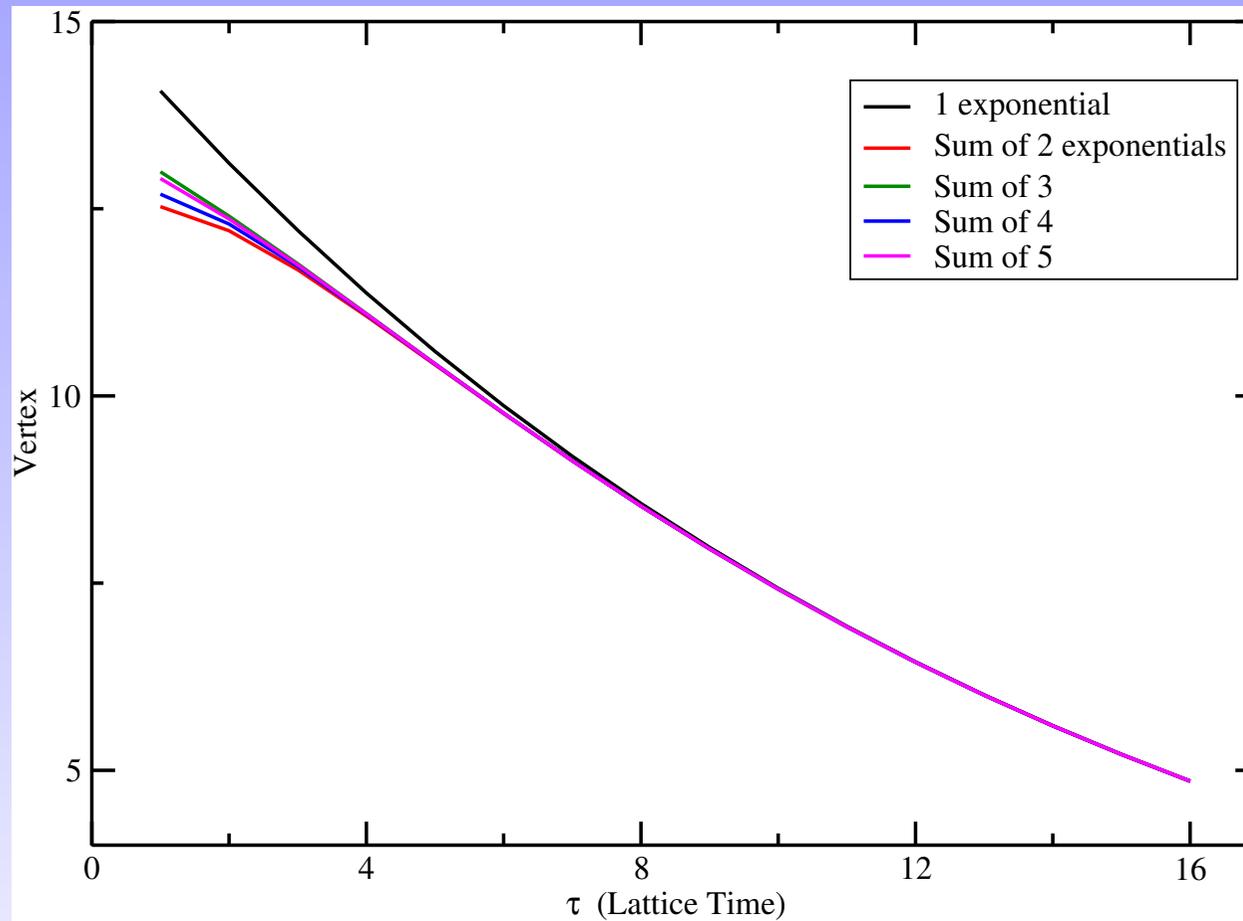
- Momentum space \rightarrow Coordinate space

$$f(p) = a + \sum_i \frac{c_i}{p^2 + m_i^2}$$

Becomes...

$$C(\tau) = a\delta(\tau) + \sum_i \frac{c_i}{2m_i} e^{-m_i\tau}$$

Lattice QCD



Effective Mass

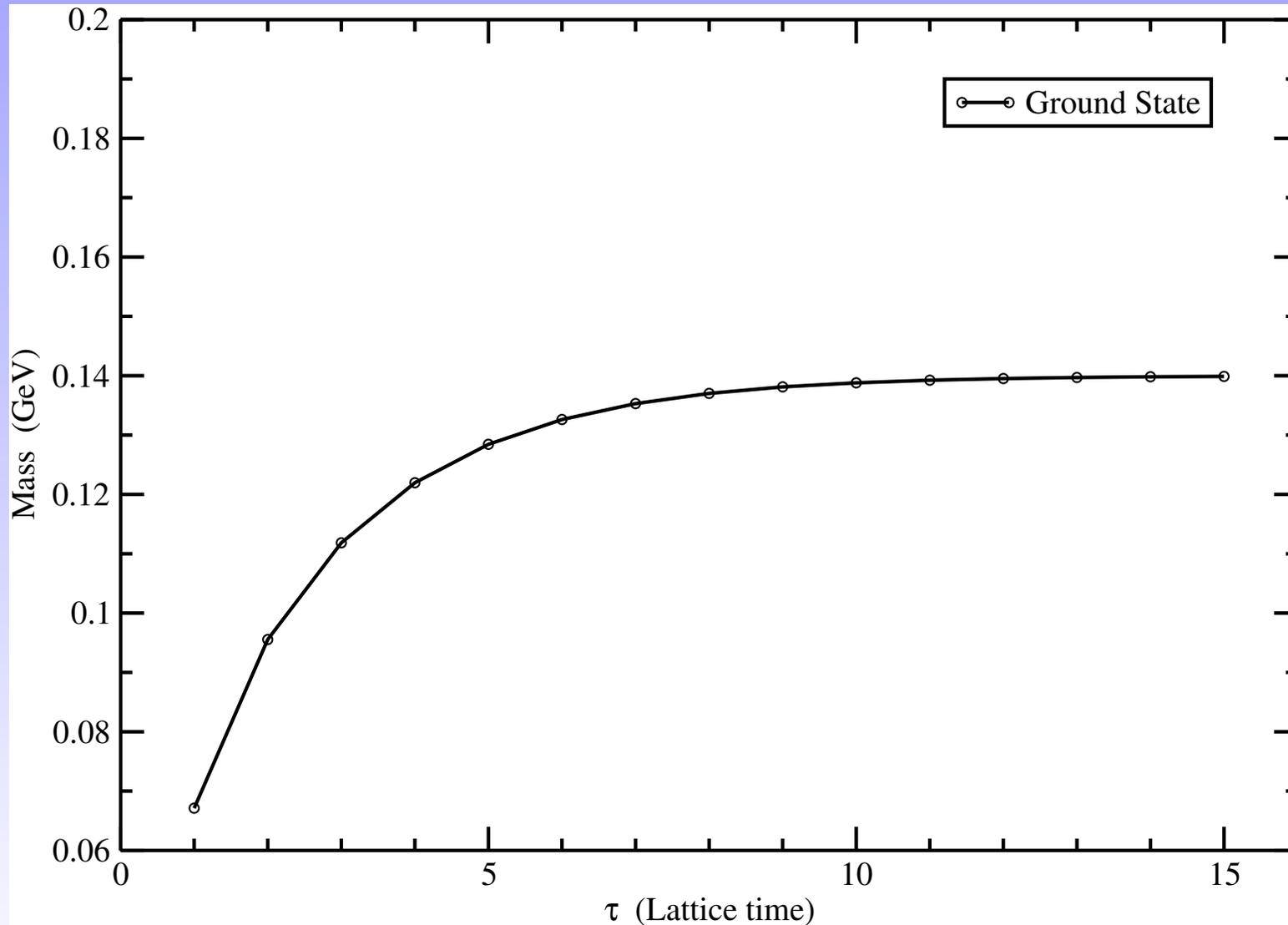
Simple to extract the ground state effective mass
ANALYTICALLY on the lattice.

$$m = -\ln \left(\frac{C(t+1)}{C(t)} \right)$$

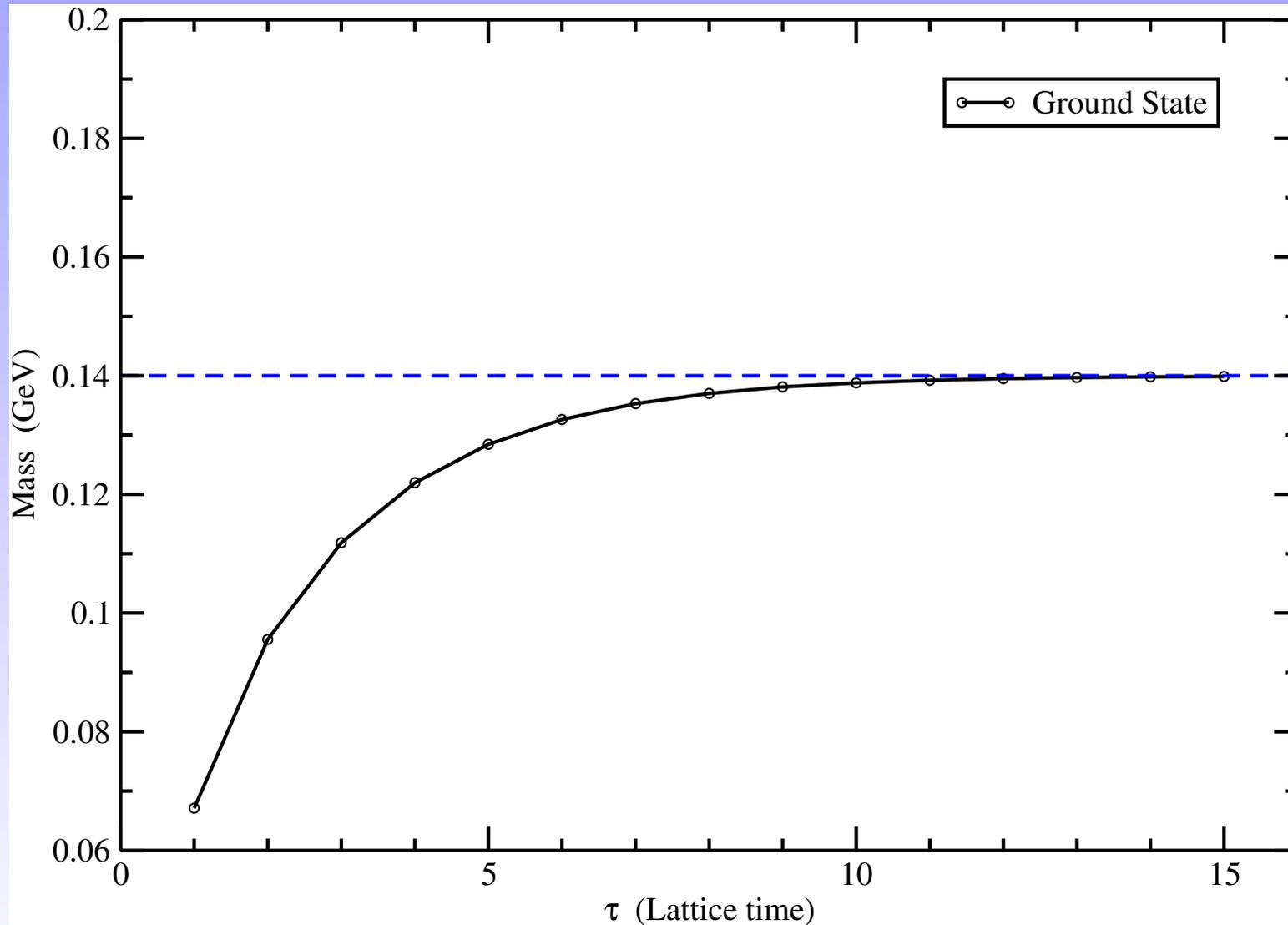
Left as an exercise for this evening.



Ground State extraction

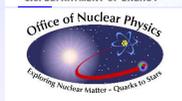


Ground State extraction

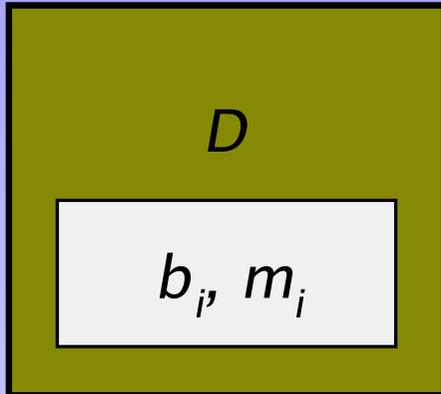


Alternate Solution

WISDOM: The best (only?) way to get access to higher excited states is through improved operators.

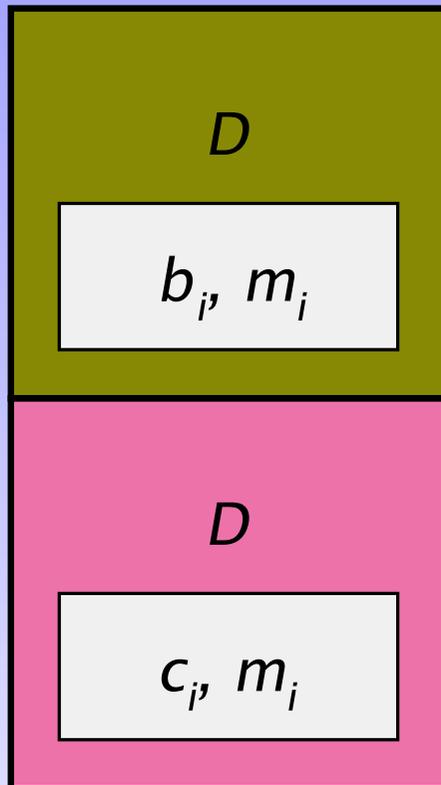


More Operators?



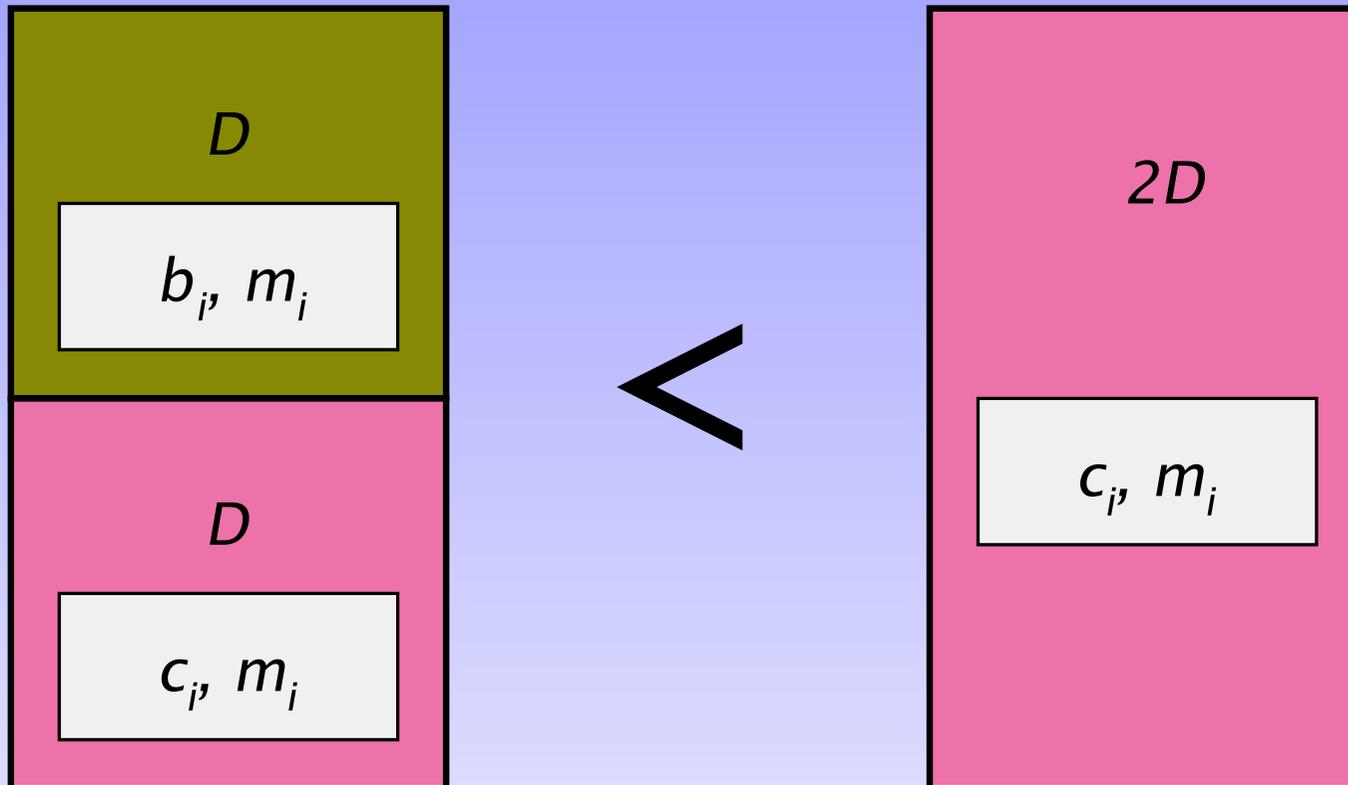
- A data set, size D gives description (b_i, m_i)

More Operators?



- A data set, size D gives description (b_i, m_i)
- A second data set, also size D gives description (c_i, m_i)

More Operators?

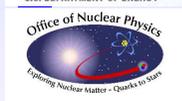


A larger data set, size $2D$ with description (c_i, m_i)
 \Rightarrow Provides **more** constraints.

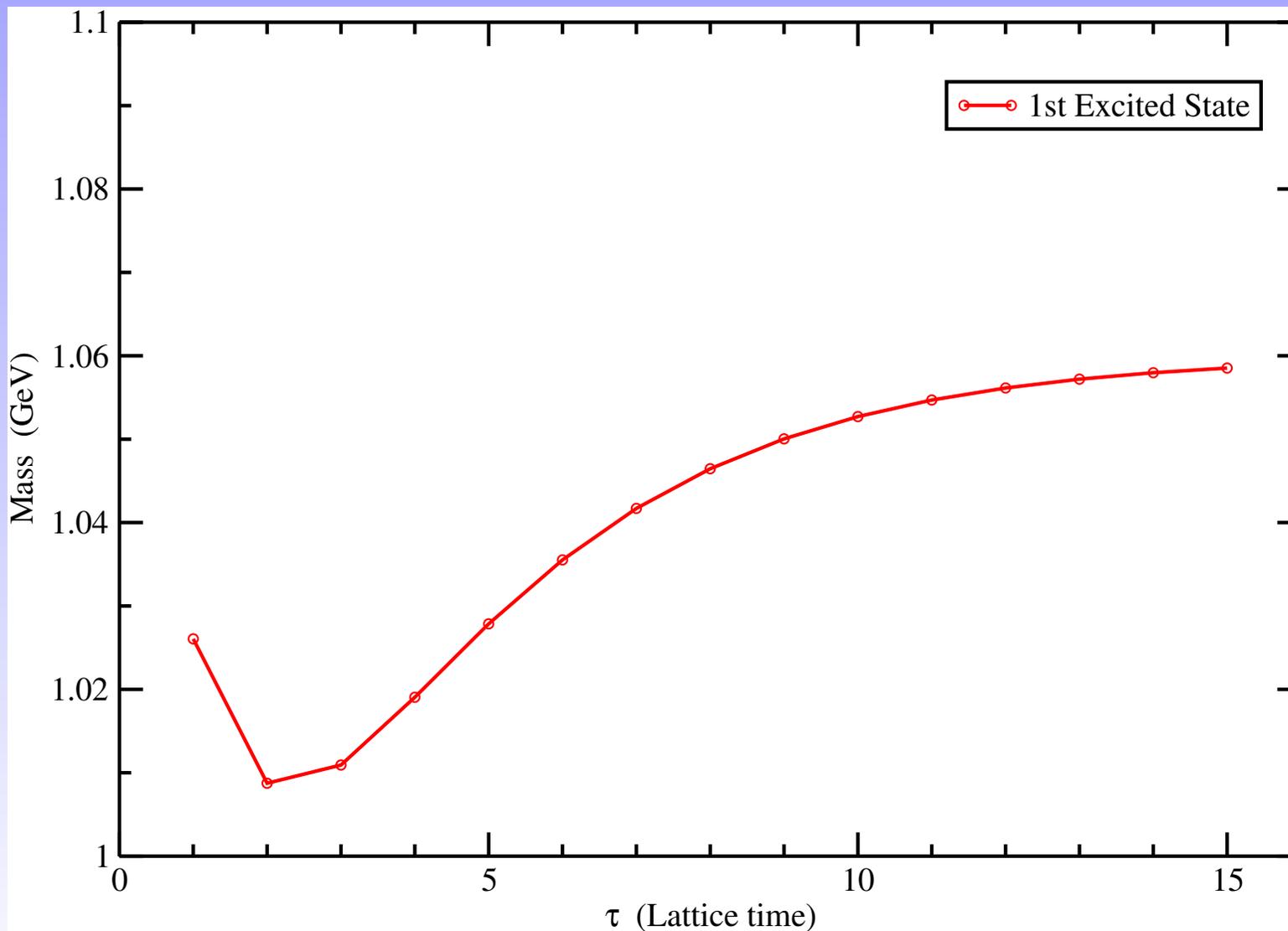
Perfect Operators

It's my model...

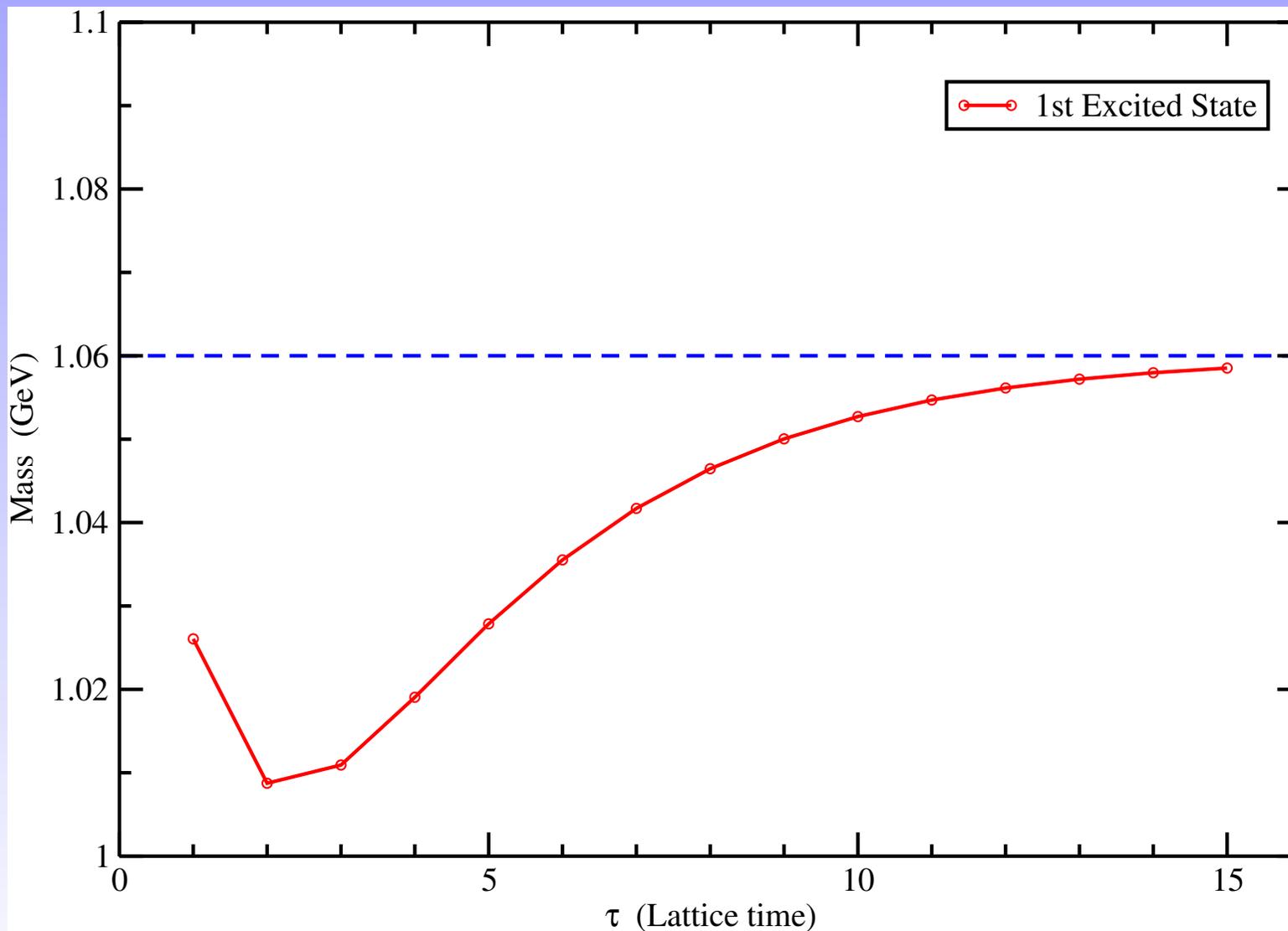
...and I'll truncate if I want to.



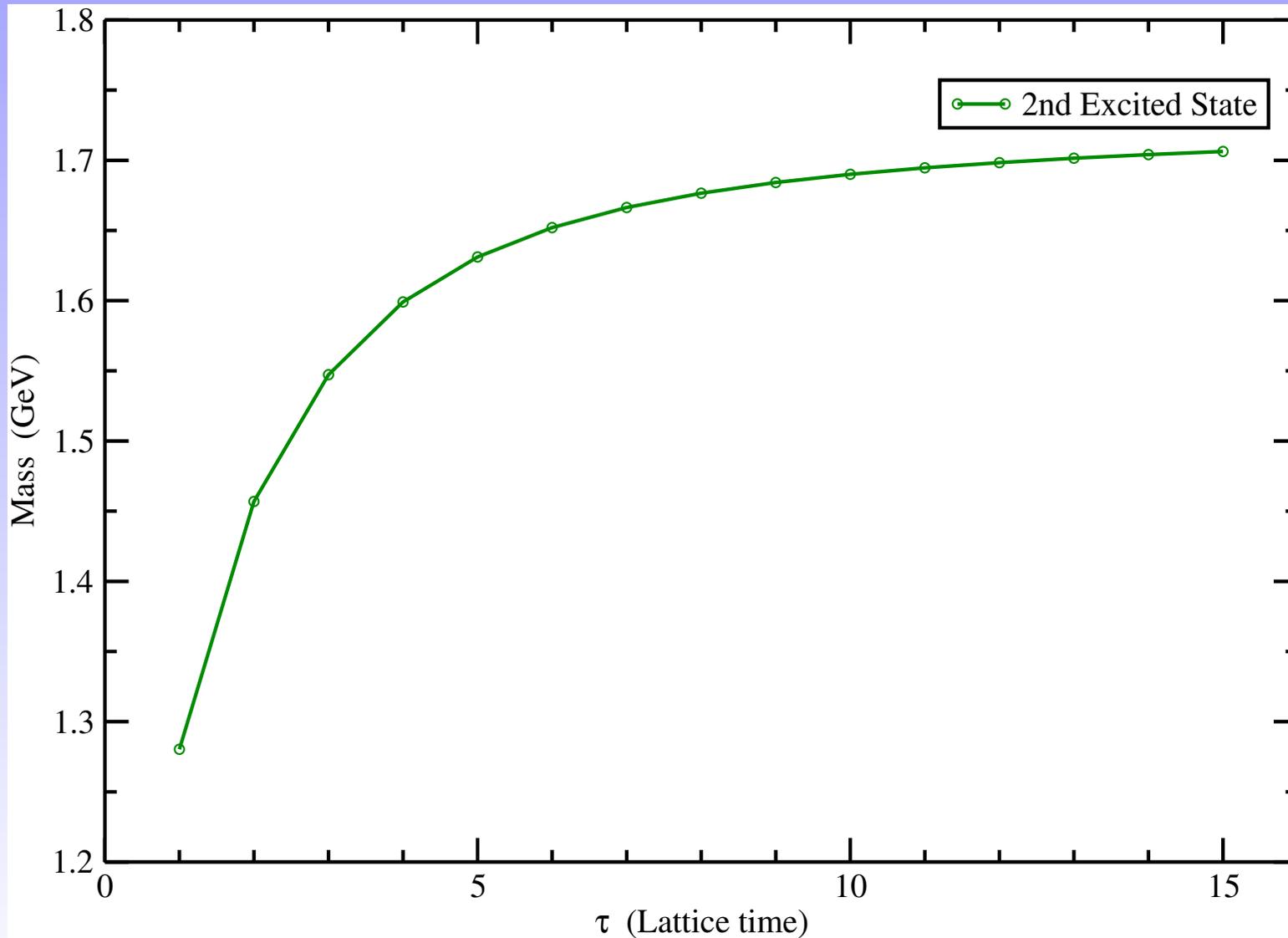
First Excited State extraction



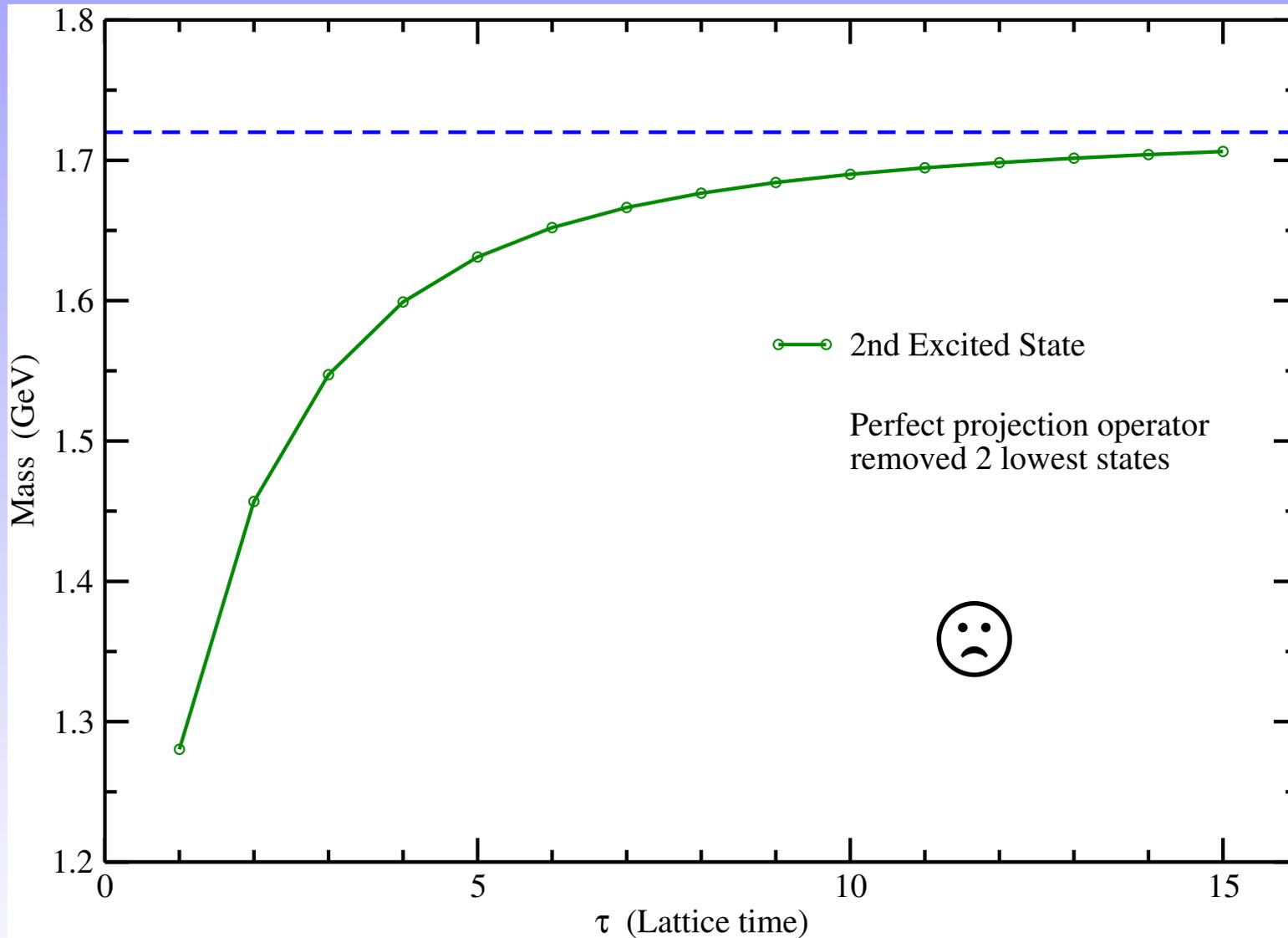
First Excited State extraction



Second Excited State extraction



Second Excited State extraction



Exponential Fits

m_0	m_1	m_2	m_3	m_4	χ^2
0.14	0.89	—	—	—	$< 10^{-06}$
0.14	1.05	1.80	—	—	$< 10^{-10}$
0.14	1.06	1.62	1.97	—	$< 10^{-11}$
0.14	1.06	1.36	1.73	1.94	$< 10^{-11}$
0.14	1.06	1.72	2.05	2.2	Source

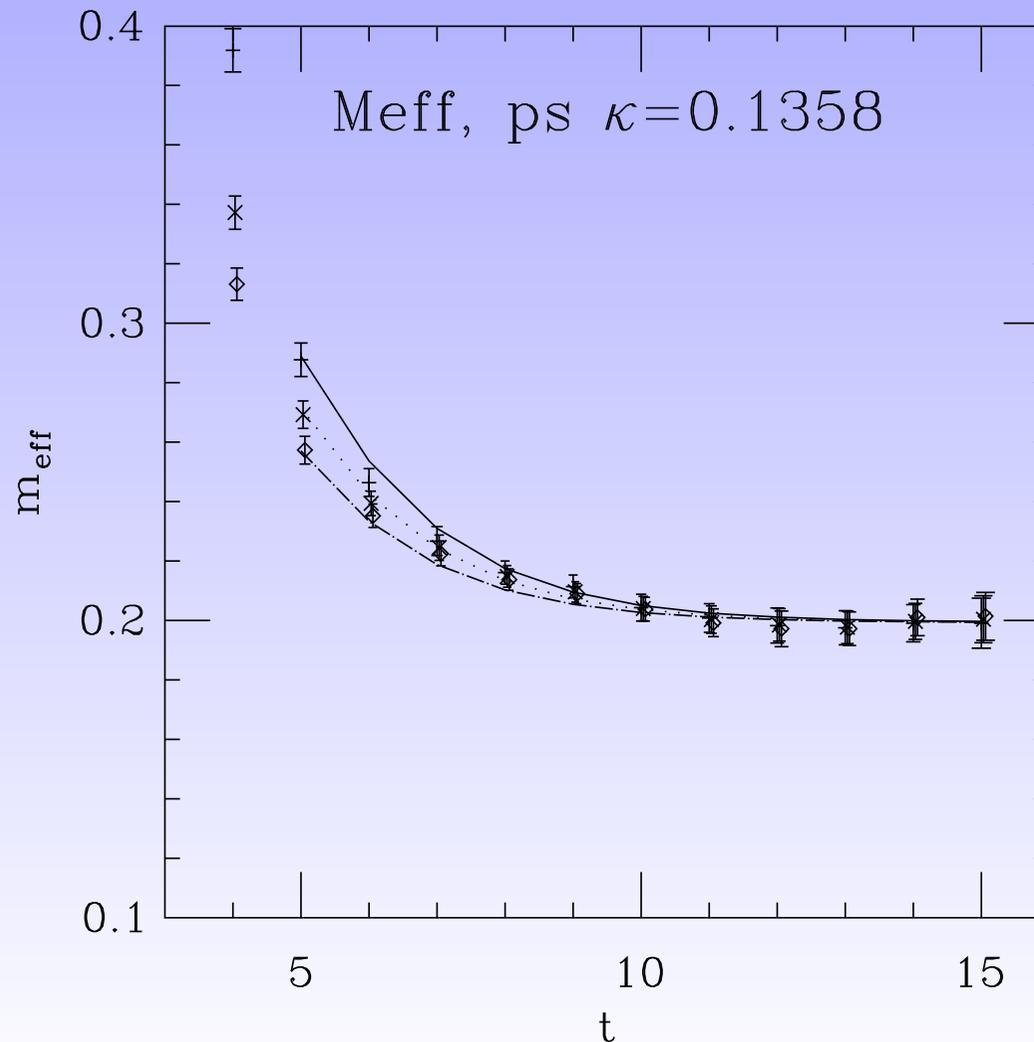
(Numbers in red are within 10%.)

(NOTE: This is χ^2 and not $\chi^2/\text{DoF}\dots$)

hep-lat/0403007

CR Allton, A Hart, D Hepburn, AC Irving, B Joo, C McNeile, C Michael, SV Wright.

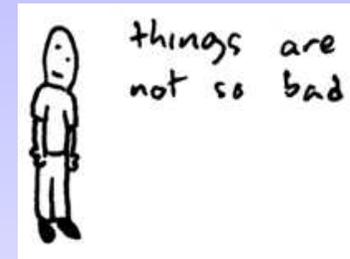
Phys.Rev. D70 (2004) 014501



Return to DSE

Strengths

- Obeys (enough) symmetries of QCD.
- FAST to calculate.
- Able to calculate in *both* Spacelike and Timelike regimes.
- Systematic errors well quantified.



Weaknesses

- Pion loops not included.
- Current implementation restricts maximum accessible meson mass.

Simplest(?) Situation

Pseudoscalar

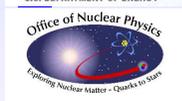
i.e. π

Light Quarks

i.e. $m_q = m_{u,d}$

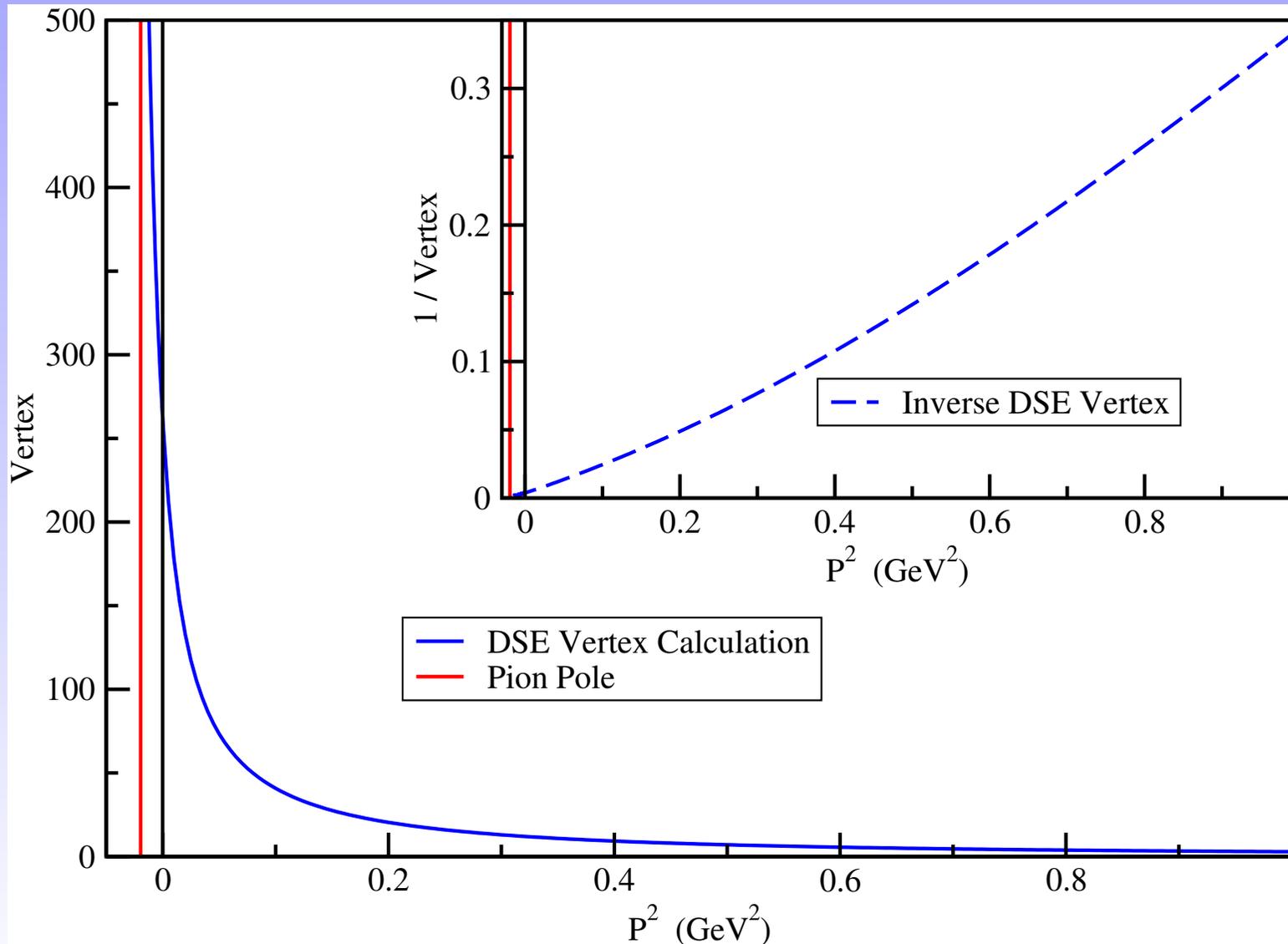
Spacelike momentum

i.e. $P^2 > 0 \text{ GeV}^2$



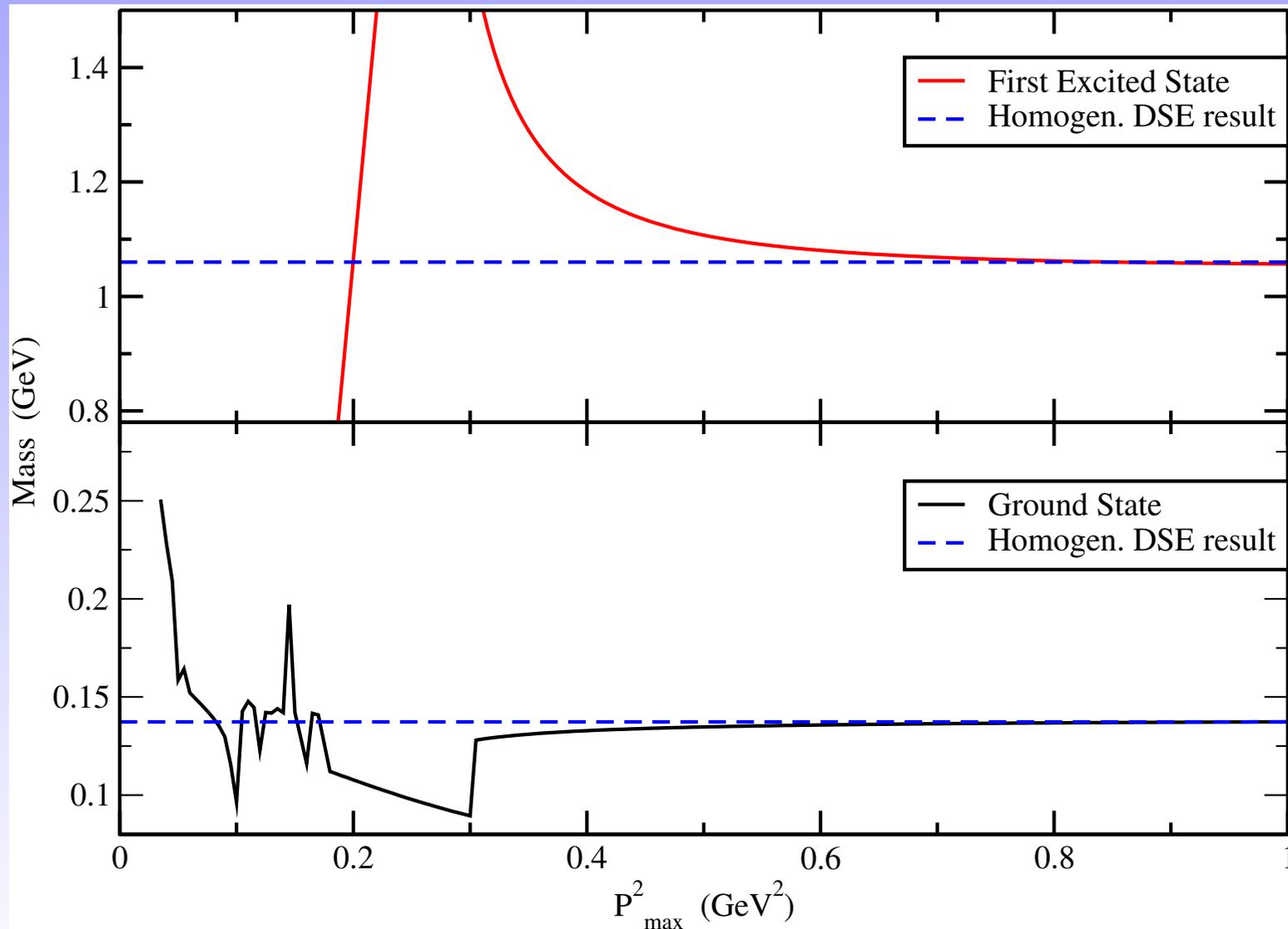
DSE Vertex Calculation

ONLY use spacelike data!



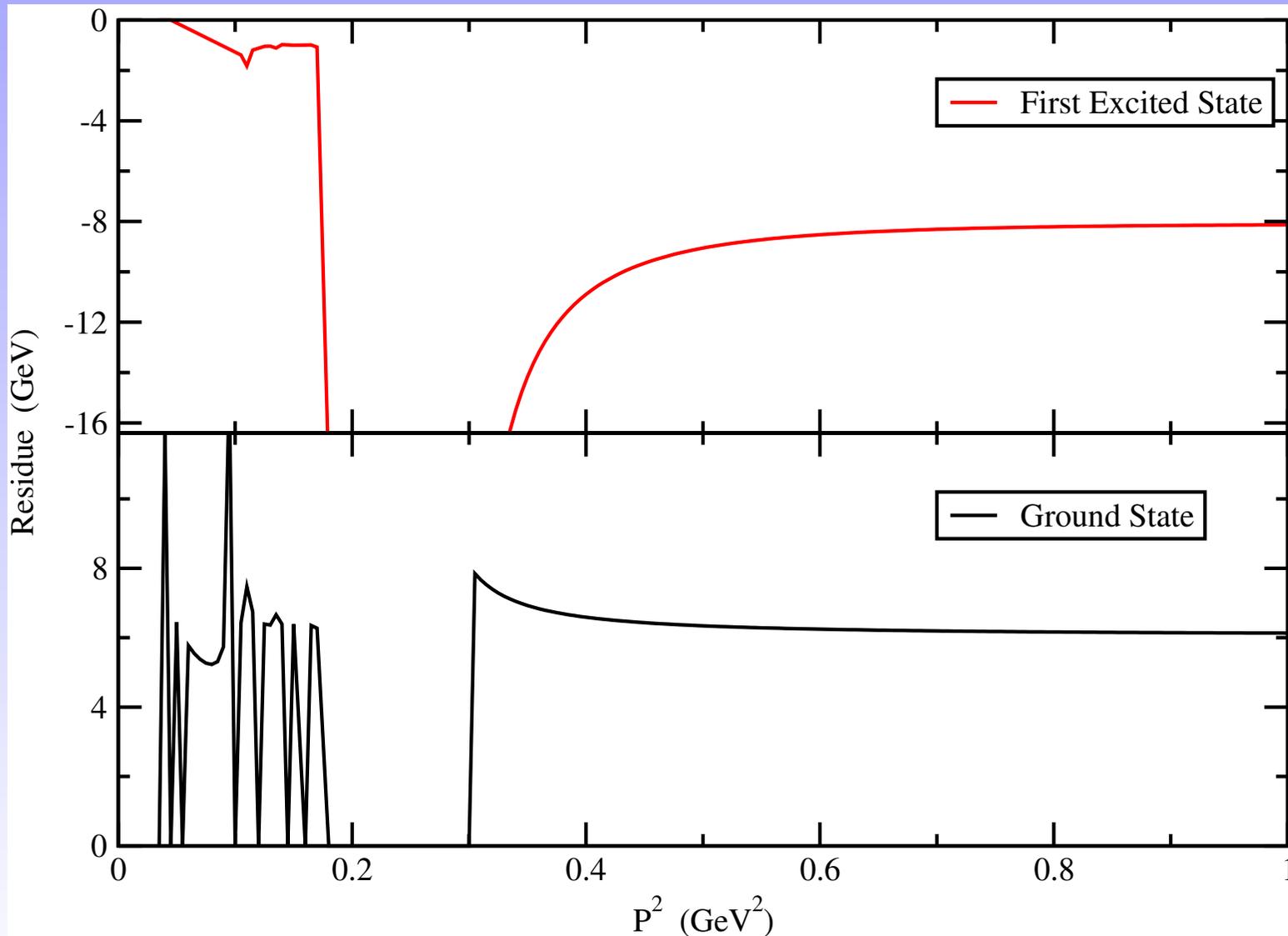
Extracted Masses

Consistent with homogeneous solution



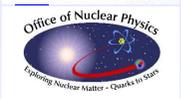
DSE Extracted Residue

Definitive evidence for sign change in residues



In a perfect world:

- Only need to know the function (and all derivatives) at a single point to uniquely determine it.



In a perfect world:

- Only need to know the function (and all derivatives) at a single point to uniquely determine it.

In the real world:

- We have FINITE PRECISION calculations.
- Cannot guarantee that there exists a unique solution of a fit of exponentials (or a Padé) to the data to arbitrary precision for:
 - Masses above ~ 1 GeV, OR
 - Second and higher excited states.



Conclusion

- Even without timelike information the extraction of the ground and first excited states is **RELIABLE**.
- The sign change in the decay constant is **REPRODUCED** *without* biasing the fit.
- With inclusion of timelike data higher excited states become **ACCESSIBLE**.

